

# Particle-in-cell simulations for Nanoplasmonic Laser Induced Fusion Experiments

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(part of NAPLIFE Collaboration)



FIAS Frankfurt Institute  
for Advanced Studies



# Nanoplasmonic Laser Inertial Fusion Experiment



Kőszeg, September 14, 2019 - Int. Workshop on Collectivity  
First meeting on the NAPLIFE project (12 people)

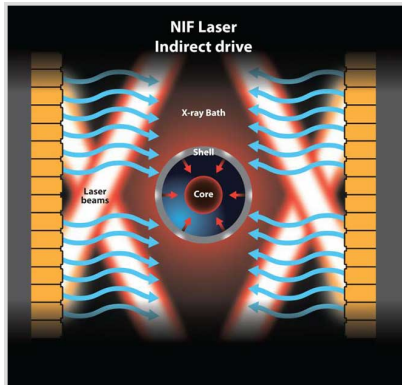
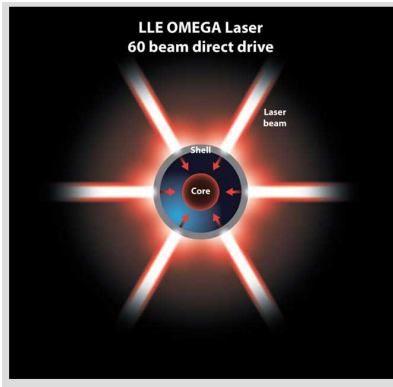
# Nanoplasmonic Laser Inertial Fusion Experiment



# Conventional Thermonuclear Fusion

- Fusion does not happen spontaneously on Earth
- Total fusion energy  $E_f = \frac{1}{4}n^2\tau\epsilon\langle v\sigma\rangle$
- $\eta E_f$  is the usable energy
- The loss is  $(1 - \eta)(E_0 + E_b)$
- $E_0 = 3nkT$ ,  $E_b = bn^2\tau\sqrt{T}$  (thermal bremsstrahlung)
- Giving the gain factor:  $Q = \frac{\eta\epsilon n\tau v\sigma}{4(1-\eta)(3kT+bn\tau\sqrt{T})}$
- $Q$  must be  $Q > 1$  for energy production
- This also means  $n\tau > \frac{3kT(1-\eta)}{\frac{1}{4}\epsilon\eta\langle v\sigma\rangle - b(1-\eta)\sqrt{T}} \rightarrow \text{LC}$
- Fulfilling the Lawson criterion
  - Magnetically confined plasmas: increase confinement time
  - Inertial confinement fusion: increase density of fusion plasma

## Direct vs Indirect drive



# Rayleigh-Taylor instabilities



# RFD

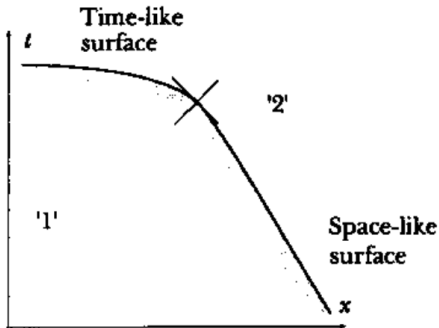


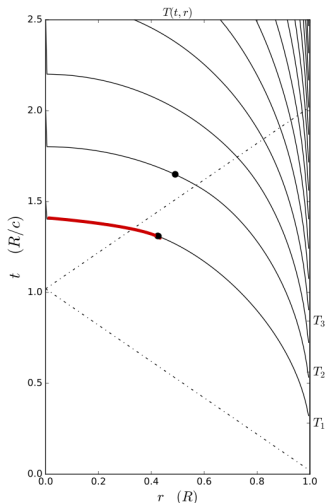
Figure 5.10: Smooth change from spacelike to timelike detonation  
[Csernai, L.P. (1987). Detonation on a time-like front for relativistic systems. Zh. Eksp. Teor. Fiz. 92, 379-386.]

# Constant absorptivity

[L.P. Csernai & D.D. Strottman, *Laser and Particle Beams* 33, 279 (2015)]

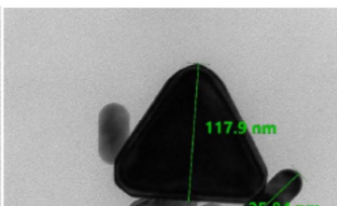
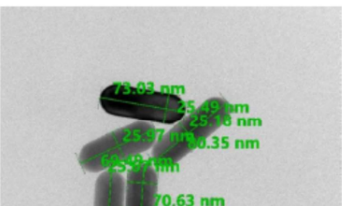
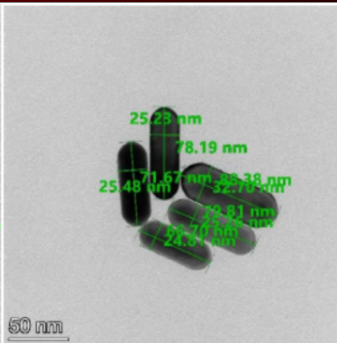
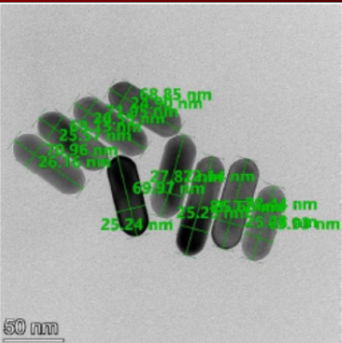
$$\alpha_{k_{middle}} = \alpha_{k_{edge}}$$

Simultaneous volume ignition is only up to 12%





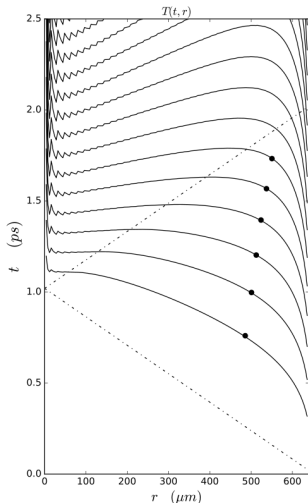
# Nanoplasmonic Laser Fusion Research Laboratory



Transmission  
Electron-  
microscopy  
photos of  
75x25 nm  
gold nano-rod  
antennas

[**Judit Kámán,**  
**A. Bonyár** et al.  
(NAPLIFE  
Collab.), Gold  
nanorods ...,  
10th ICNFP  
2021, **Kolymbari**]

## Changing absorptivity

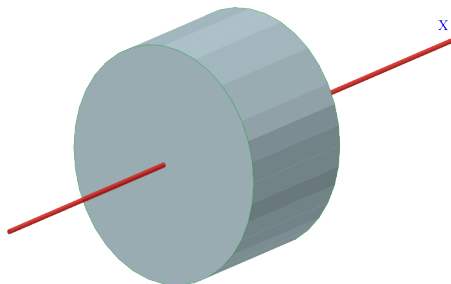


[Csernai, L.P., Kroo, N. and Papp, I. (2017). Procedure to improve the stability and efficiency of laser-fusion by nano-plasmonics method. Patent P1700278/3 of the Hungarian Intellectual Property Office.]

$$\alpha_{k_{middle}} \approx 4 \times \alpha_{k_{edge}}$$

Simultaneous volume ignition is up to 73%

# Flat target

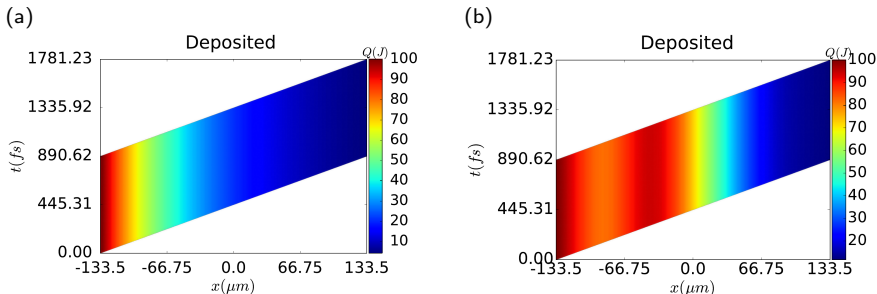


Schematic view of the cylindrical, flat target of radius,  $R$ , and thickness,  $h$ .

$$V = 2\pi R^3, \quad R = \sqrt[3]{V/(2\pi)}, \quad h = \sqrt[3]{4V/\pi}.$$

[L.P. Csernai, M. Csete, I.N. Mishustin, A. Motorenko, I. Papp, L.M. Satarov, H. Stöcker & N. Kroó, Radiation- Dominated Implosion with Flat Target, *Physics and Wave Phenomena*, **28** (3) 187-199 (2020)]

# Varying absorptivity



**Deposited energy** per unit time in the space-time plane across the depth,  $h$ , of the flat target. **(a) without nano-shells (b) with nano-shells**  
 To increase central absorption we used the following distribution:

$$\alpha_{ns}(s) = \alpha_{ns}^C + \alpha_{ns}(0) \cdot \exp\left[\frac{4x^2}{L^2 - x^2}\right].$$

# Particle In Cell methods

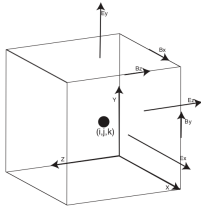


Figure 1. Yee staggered grid used for the Maxwell solver in EPOCH.

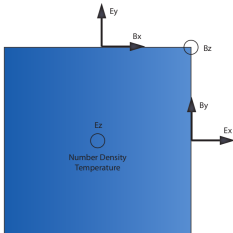


Figure 2: The Yee grid in 2D

[F.H. Harlow (1955). A Machine Calculation Method for Hydrodynamic Problems. Los Alamos Scientific Laboratory report LAMS-1956]

[T.D. Arber et al 2015 Plasma Phys. Control. Fusion 57 113001]

A **super-particle** (**marker-particle**) is a computational particle that represents many real particles.

Particle **mover** or **pusher** algorithm as standard **Boris algorithm**.

**Finite-difference time-domain method** for solving the time evolution of **Maxwell's equations**.

# FDTD in EPOCH

- $\mathbf{E}_{n+\frac{1}{2}} = \mathbf{E}_n + \frac{\Delta t}{2} \left( c^2 \nabla \times \mathbf{B}_n - \frac{\mathbf{j}_n}{\epsilon_0} \right)$
- $\mathbf{B}_{n+\frac{1}{2}} = \mathbf{B}_n - \frac{\Delta t}{2} \left( \nabla \times \mathbf{E}_{n+\frac{1}{2}} \right)$
- Call particle pusher which calculates  $\mathbf{j}_{n+1}$
- $\mathbf{B}_{n+1} = \mathbf{B}_{n+\frac{1}{2}} - \frac{\Delta t}{2} \left( \nabla \times \mathbf{E}_{n+\frac{1}{2}} \right)$
- $\mathbf{E}_{n+1} = \mathbf{E}_{n+\frac{1}{2}} + \frac{\Delta t}{2} \left( c^2 \nabla \times \mathbf{B}_{n+1} - \frac{\mathbf{j}_{n+1}}{\epsilon_0} \right)$

# Particle pusher

- Solves the relativistic equation of motion under the Lorentz force for each marker-particle

$$\mathbf{p}_{n+1} = \mathbf{p}_n + q\Delta t \left[ \mathbf{E}_{n+\frac{1}{2}}(\mathbf{x}_{n+\frac{1}{2}}) + \mathbf{v}_{n+\frac{1}{2}} \times \mathbf{B}_{n+\frac{1}{2}}(\mathbf{x}_{n+\frac{1}{2}}) \right]$$

$\mathbf{p}$  is the particle momentum  $q$  is the particle's charge  $\mathbf{v}$  is the velocity.

$\mathbf{p} = \gamma m \mathbf{v}$ , where  $m$  is the rest mass  $\gamma = [(\mathbf{p}/mc)^2 + 1]^{1/2}$

- Villasenor and Buneman current deposition scheme [Villasenor J & Buneman O 1992 Comput. Phys. Commun. 69 306], always satisfied:  $\nabla \cdot \mathbf{E} = \rho/\epsilon_0$ , where  $\rho$  is the charge density.

## Particle shape

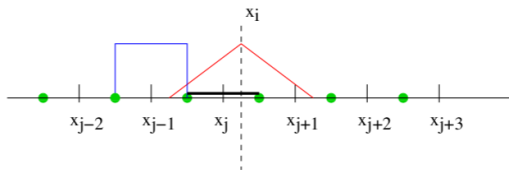


Figure 3: Second order particle shape function

First order approximations are considered

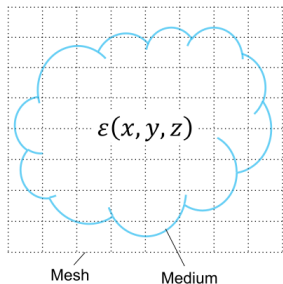
$$F_{part} = \frac{1}{2} F_{i-1} \left( \frac{1}{2} + \frac{x_i - X}{\Delta x} \right)^2 + \frac{1}{2} F_i \left( \frac{3}{4} - \frac{(x_i - X)^2}{\Delta x^2} \right)^2 + \frac{1}{2} F_{i+1} \left( \frac{1}{2} + \frac{x_i - X}{\Delta x} \right)^2$$

[EPOCH 4.0 dev manual]

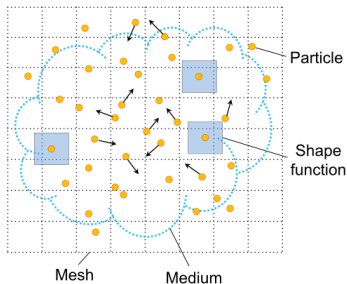


# Nanorod

A Field simulation



B Particle simulation



[W. J. Ding, et al., Particle simulation of plasmons Nanophotonics, vol. 9, no. 10, pp. 3303-3313 (2020)]

# Nanorod

## Typical Field solver:

$$\epsilon(\omega) = 1 - \frac{\omega_p^2}{(\omega^2 + i\gamma\omega)}$$

where  $\omega_p$  is the plasma frequency:  $\sqrt{\frac{n_e e^2}{m' \epsilon_0}}$

$\gamma$  is the damping factor or collision frequency:  $\gamma = \frac{1}{\tau}$  and  $\tau$  is the average time between collisions

## Particle simulation:

$$\frac{\partial \mathbf{E}}{\partial t} = \frac{1}{\mu_0 \epsilon_0} \nabla \times \mathbf{B} - \frac{\mathbf{J}}{\epsilon_0}, \quad \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

$\gamma_i m_i \mathbf{v}_i = q_i (\mathbf{E}_i + \mathbf{v}_i \times \mathbf{B}_i)$ ,  $\gamma_i$  is the relativistic factor

# Metal Nanoparticles as Plasmas

The conduction band electrons in metals behave as strongly coupled plasmas.

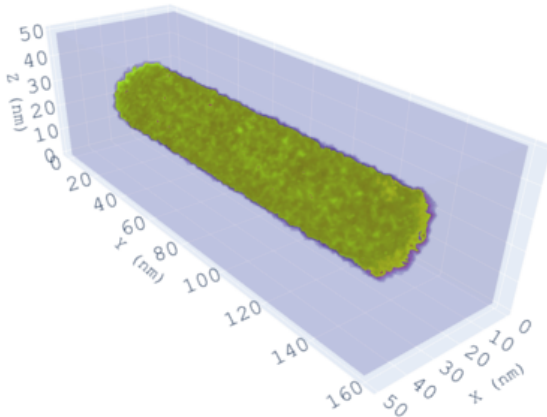
For golden nanorods of 25nm diameter in vacuum this gives an effective wavelength of  $\lambda_{eff} = 266\text{nm}$

$$\frac{\lambda_{eff}}{2R\pi} = 13.74 - 0.12[\epsilon_{\infty} + 141.04] - \frac{2}{\pi} + \frac{\lambda}{\lambda_p} 0.12\sqrt{\epsilon_{\infty} + 141.04}$$

[Lukas Novotny, Effective Wavelength Scaling for Optical Antennas, Phys. Rev. Lett. **98**, 266802 (2007).]

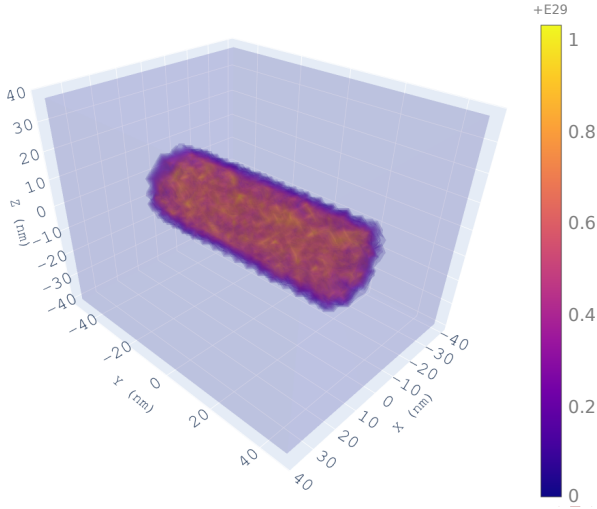
# Kinetic Modelling of the Nanorod

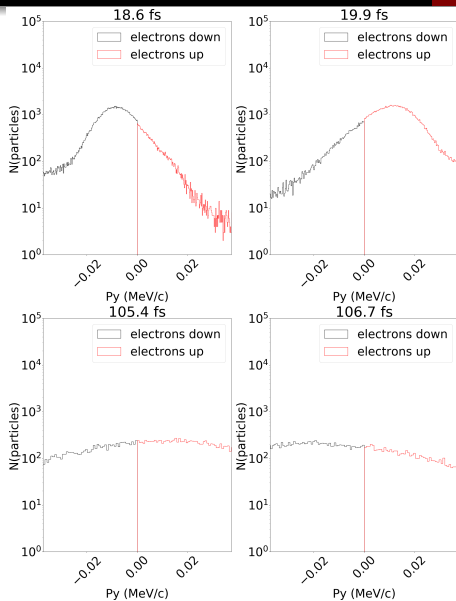
Nanorod inside a PIC simulation box in vacuum



# Kinetic Modelling of the Nanorod

Nanorod inside a PIC simulation box in UDMA





Considerations for the simulation box:

$$S_{CB} = 530 \times 530 \text{ nm}^2 =$$

$$2.81 \times 10^{-9} \text{ cm}^2 \text{ and length of}$$

$$L_{CB} = 795 \text{ nm}$$

beam crosses the box in

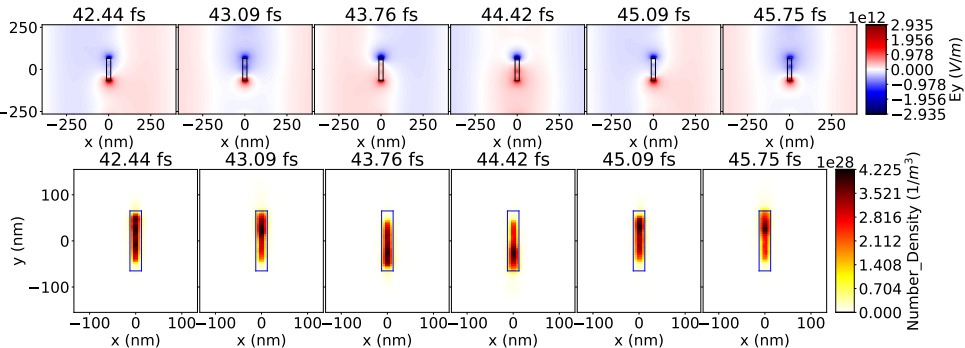
$$T = 795 \text{ nm}/c = 2.65 \text{ fs}$$

Nanorod size: **25 nm** diameter  
 with **130 nm** length

Pulse length:  $40 \times \lambda/c = 106 \text{ fs}$

Intensity:  $4 \times 10^{15} \text{ W/cm}^2$

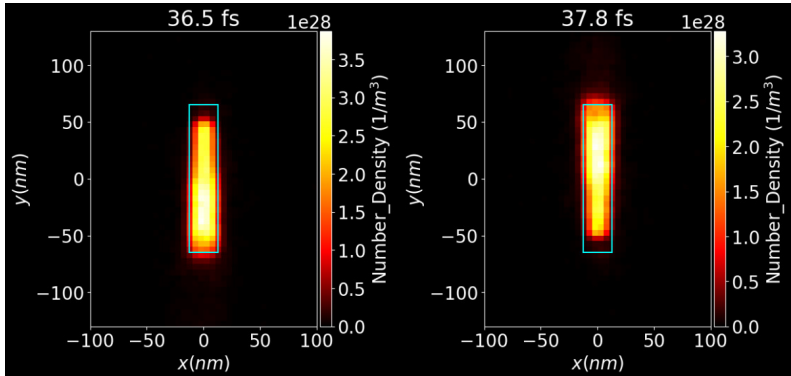
# Kinetic Modelling of the Nanorod



- Evolution of the  $E$  field's  $y$  component from 42.4 till 45.7 fs, around a nanorod of 25x130 nm.
- The direction of the  $E$  field at the two ends of the nanorod does not change.

# Kinetic Modelling of the Nanorod

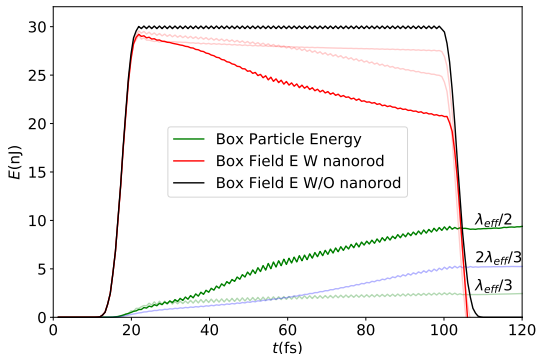
Evolution of the nanoantenna



Number density of electrons in the middle of a nanorod of size 25x130 nm at different times. The nanorod is orthogonal to the beam direction,  $x$ .



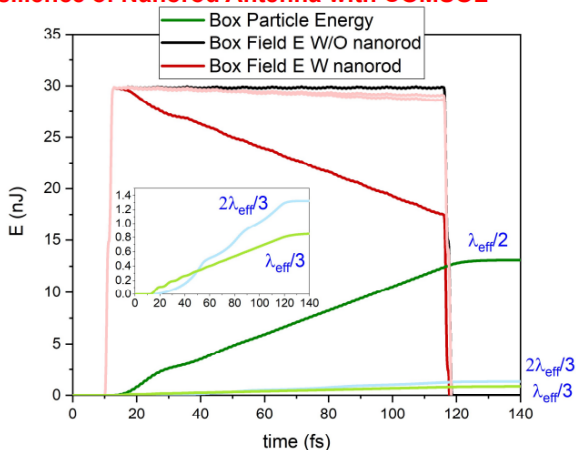
## In vacuum



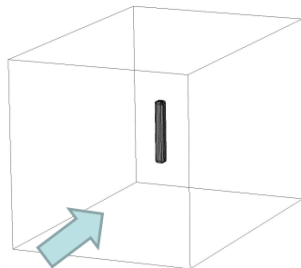
energy in the box **without nanorod** antenna  $3 \times 10^{-8}$  J (black line)  
**nanorod** absorbs EM energy reducing it to  $2.3 \times 10^{-8}$  J (red line)  
**deposited** energy in the nanorod (green line)  
results in light absorption cross section highest

# Comparison with other methods (Csernai, Csete et al.)

## Resilience of Nanorod Antenna with COMSOL

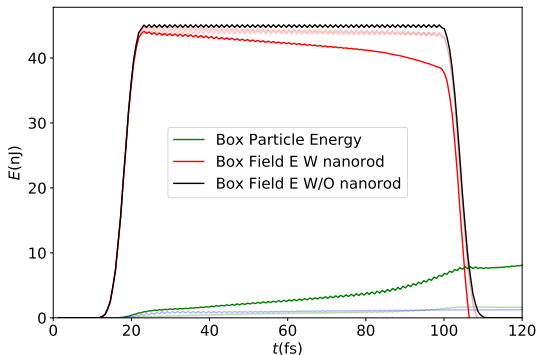


FEM computations with the same model parameters



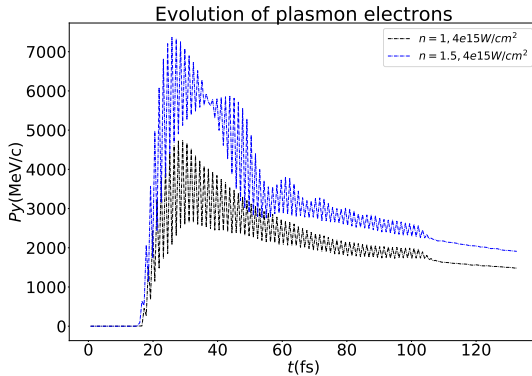
Good qualitative agreement between FEM and EPOCH/PIC methods  
 Quantitative difference:

# In UDMA

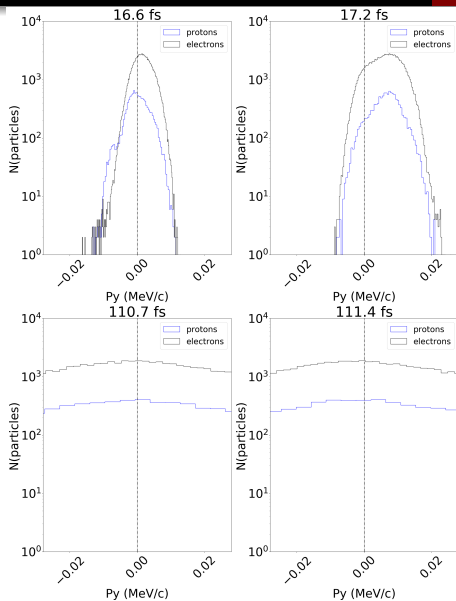


deposited energy in the nanorod (green line)

# In UDMA and vacuum



accumulated momentum of conduction electrons in **vacuum** (blue) and in **UDMA** (black) with their corresponding resonant length



## Protons surrounding the nanorod

Considerations for the simulation box:

$S_{CB} = 530 \times 530 \text{ nm}^2 = 2.81 \times 10^{-9} \text{ cm}^2$  and length of  $L_{CB} = 795 \text{ nm}$

beam crosses the box in  $T = 795 \text{ nm}/c = 2.65 \text{ fs}$

Nanorod size: **25 nm** diameter with **85 nm** length

Pulse length:  $40 \times \lambda/c = 106 \text{ fs}$   
 Intensity:  $4 \times 10^{15} \text{ W/cm}^2$

# Conclusion and Outlook

- Our results agree with the those of Mária Csete in vacuum
- Quantitative differences mainly come at different lengths from resonance
- Levitation effect comes only in vacuum, needs further investigation
- Next step is estimating the target pre-compression