

Femtoscscopy for the NAPLIFE nano-fusion project?

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on Multiparticle Dynamics,
Aug. 24, 2023, Gyöngyös, Hungary

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NAPLIFE Collaboration
Univ. of Bergen, Norway

How to remedy the problems of present Laser Fusion trials of NIF@Livermore & OMEGA@Rochester

Two ideas are combined by L.P. Csernai, N. Kroo, I. Papp:
[Patent # P1700278/3] (2017)

Problems:

- Rayleigh-Taylor instability
- Slow propagation of burning from central hot-spot

Solution:

- **Heat the system uniformly by radiation with RFD (1)**
- **Achieve uniform heating by Nano-Technology (2)**

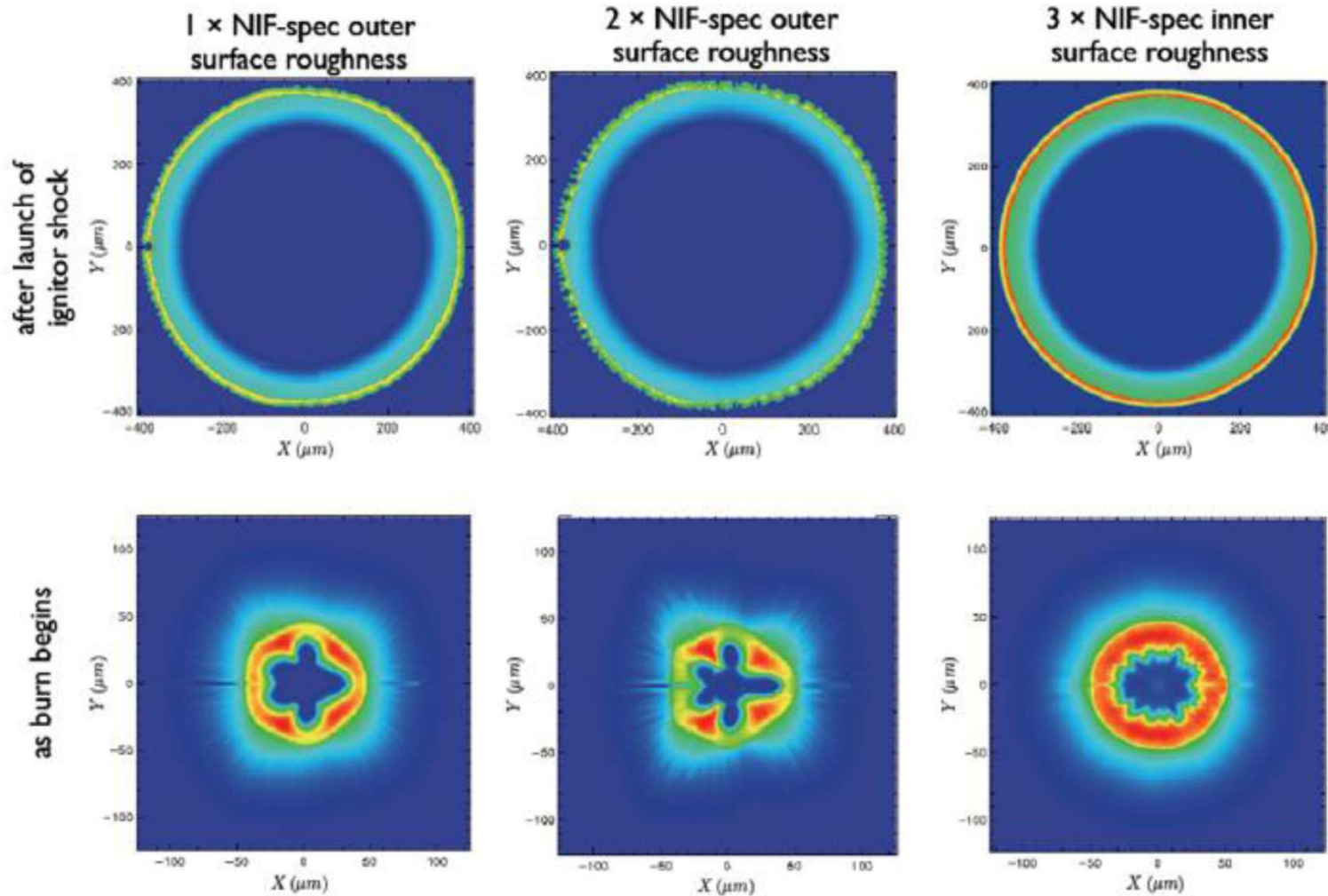
[L.P. Csernai, N. Kroo, I. Papp, *Laser and Particle Beams*, LPB, 36(2), (2018) 171-178. .

<https://doi.org/10.1017/S0263034618000149>]

But let us go back in history →

Rayleigh-Taylor Instability

NIF – RT instability



The target is compressed to density $\sim 700 \text{ g/cm}^3$.

But, although an ablator layer is used, only $\sim 10\%$ -of the target is ignited.

Elsewhere the surface protruded as “potato from the potato press”:
RT- instability.

How can we prevent it

Idea - #1

[A.H. Taub (1948)]

PHYSICAL REVIEW

VOLUME 74, NUMBER 3

AUGUST 1, 1948

Relativistic Rankine-Hugoniot Equations

A. H. TAUB

*University of Illinois, Urbana, Illinois and Institute for Advanced Study, Princeton University, Princeton, New Jersey**

Next we suppose that the three-dimensional volume is a shell of thickness ϵ enclosing a surface of discontinuity Σ whose three-dimensional normal vector is Λ_i . If we choose our coordinate system so that the discontinuity is at rest, then since

$$\underline{\lambda_\alpha \lambda^\alpha = 1}, \quad \sum_{i=1}^3 \Lambda_i^2 = 1,$$

we have

$$\lambda_i = \Lambda_i \quad \text{and} \quad \underline{\lambda_4 = 0}.$$

Hence Eqs. (7.1) and (7.2) become, as ϵ goes to zero,

$$[\rho^0 u^i \Lambda_i] = 0, \quad (7.3)$$

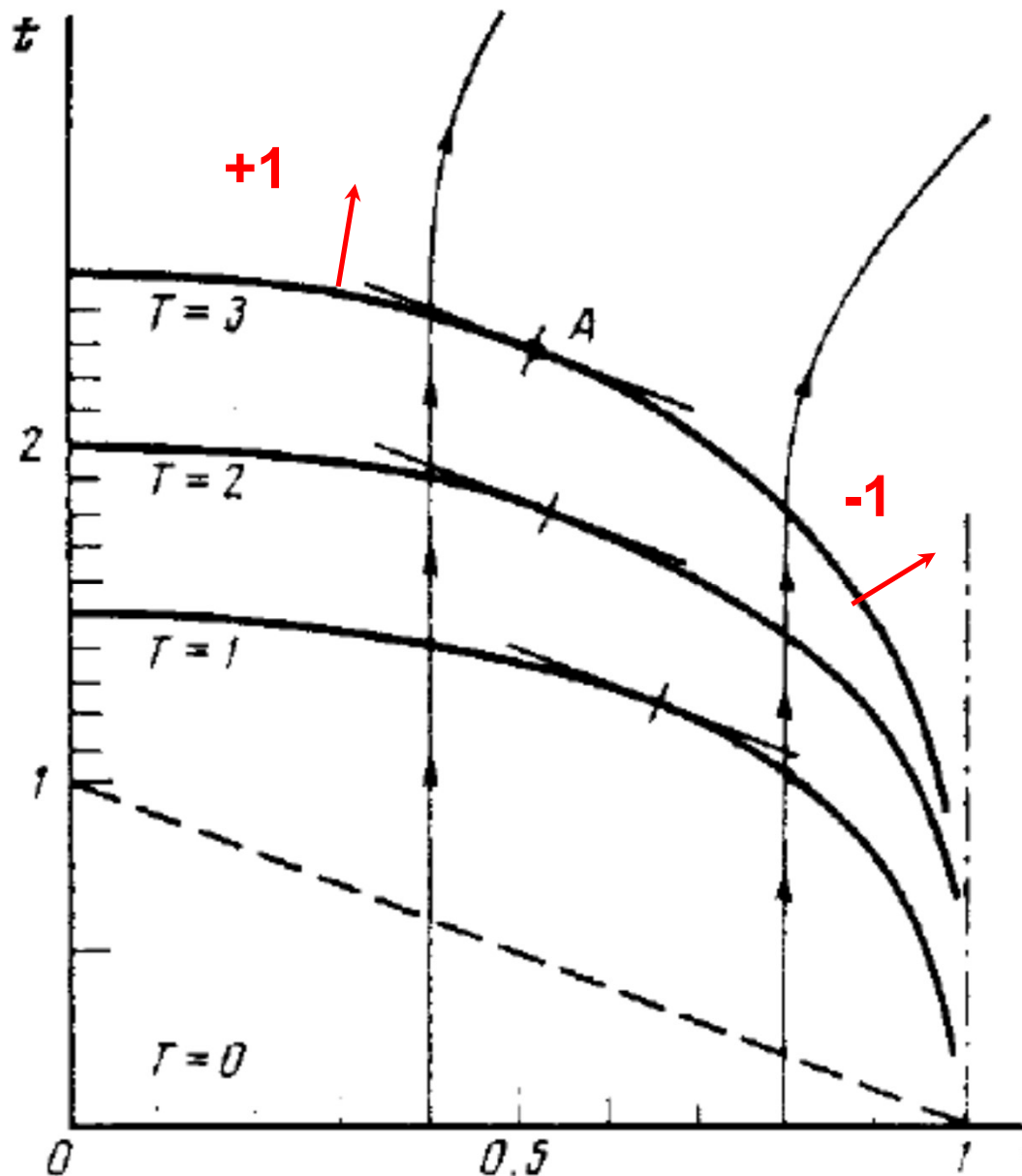
$$[T^{\alpha i} \Lambda_i] = 0, \quad (7.4)$$

where

$$[f] = f_+ - f_-$$

Taub assumed that (physically) only slow space-like shocks or discontinuities may occur (with space-like normal, $\lambda_4=0$).

This was then taken as standard, since then (e.g. LL 1954-)



[L. P. Csernai, Zh. Eksp. Teor. Fiz. 92, 379-386 (**1987**) & Sov. Phys. JETP 65, 216-220 (1987)]

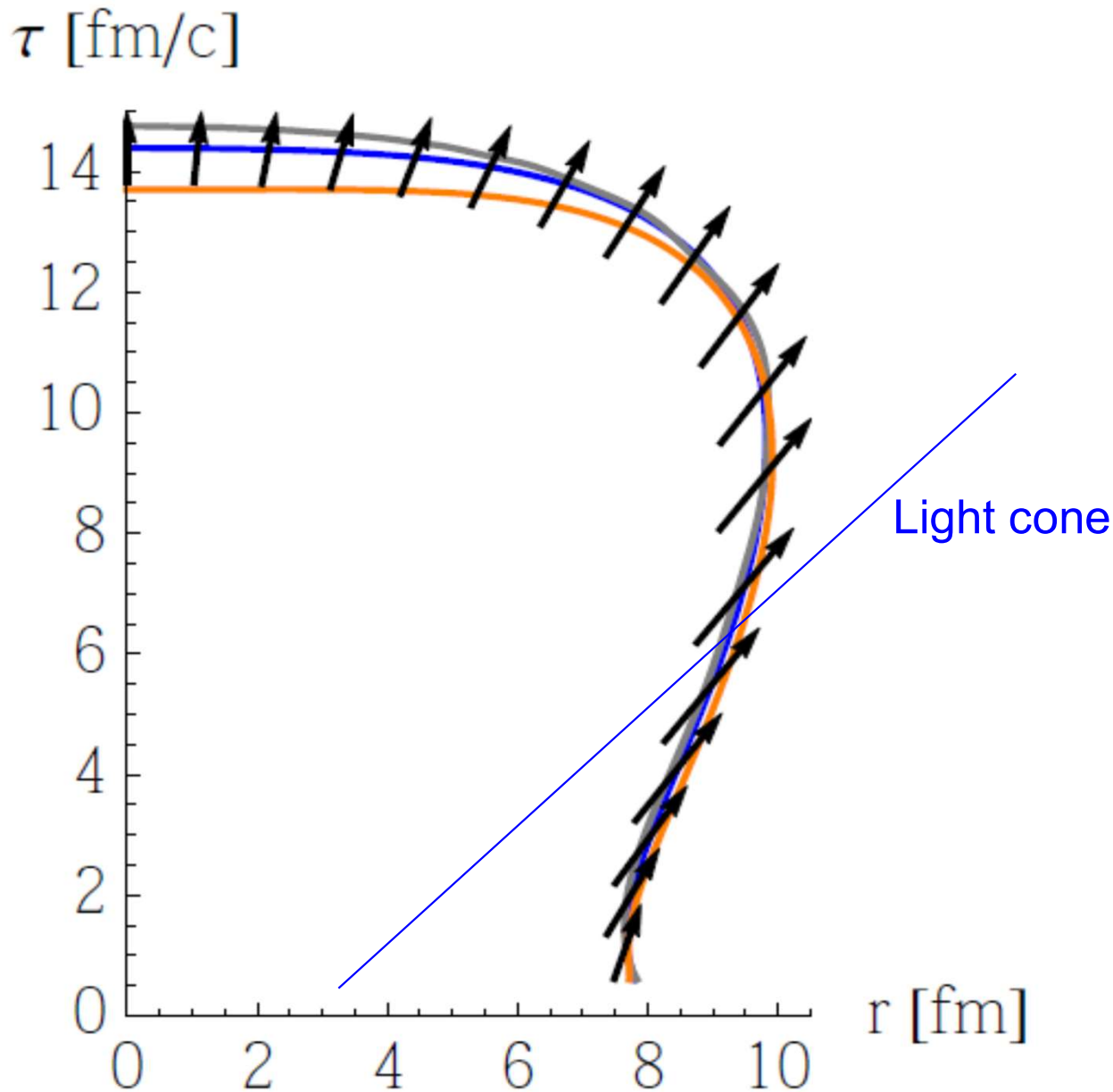
corrected the work of [**A. Taub**, Phys. Rev. 74, 328 (**1948**)]

$$\lambda_\alpha \lambda^\alpha = \pm 1$$

Л. П. Чернаи

**ДЕТОНАЦИЯ НА ВРЕМЕНИПОДОБНОМ ФРОНТЕ
ДЛЯ РЕЛЯТИВИСТСКИХ СИСТЕМ**

Журнал экспериментальной и теоретической физики



@ CERN in High energy heavy ion collisions

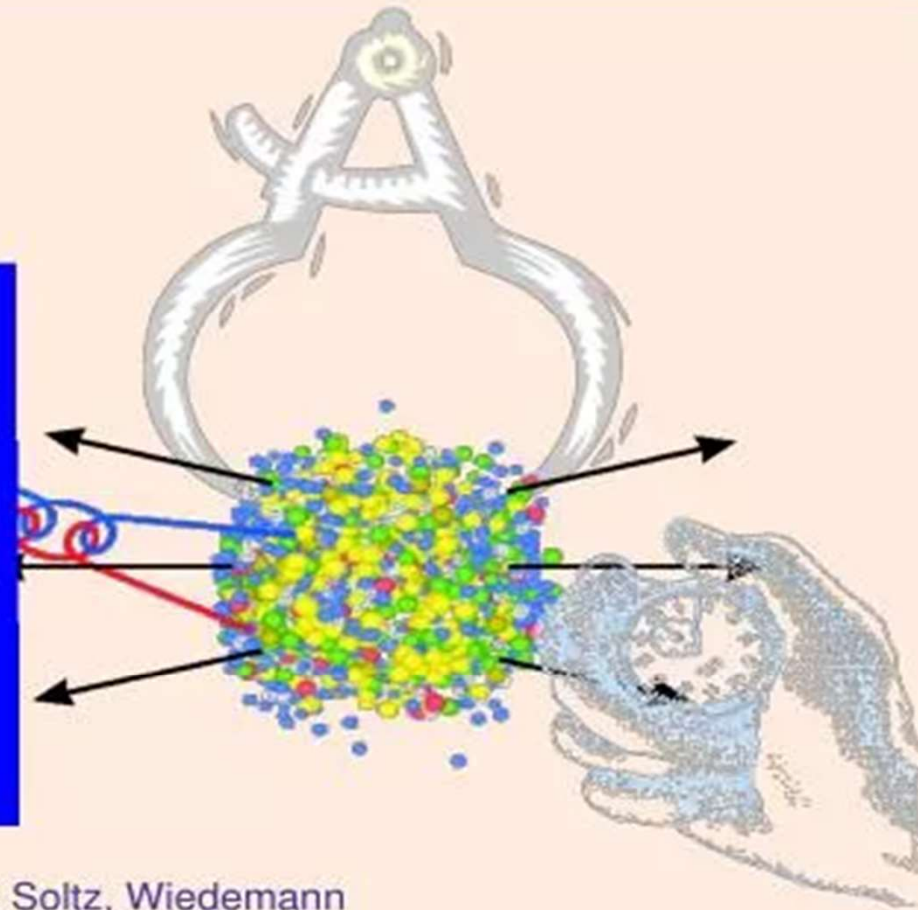
[Stefan Floerchinger, and Urs Achim Wiedemann, Phys. Rev. C 89, 034914 (2014)]

Femtoscscopy in heavy ion collisions: Wherefore, Whence, & Whither?

Mike Lisa

Ohio State University

- Wherefore (=“why?”)
 - motivation & (basic) formalism
- Whence (=“from where?”)
 - systematics over 2 decades
- Whither (=“to where?”)
 - or “wither”...?

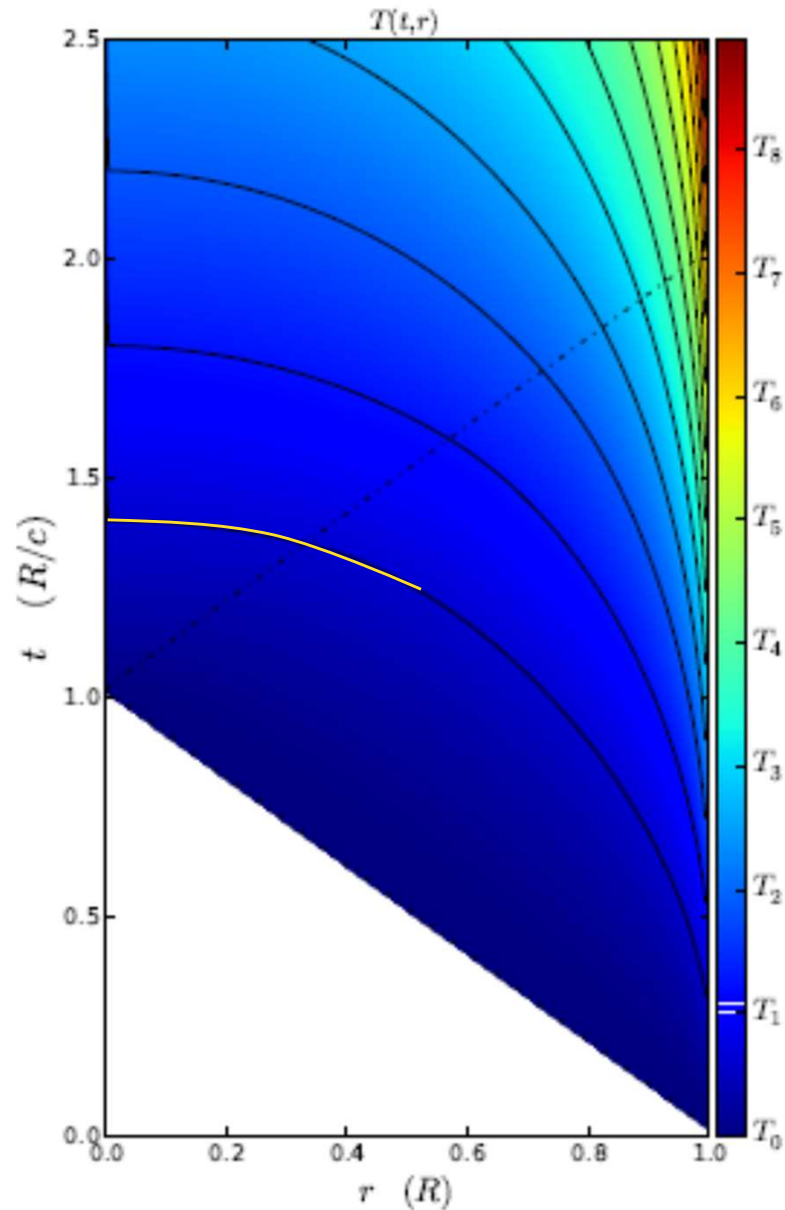


<http://www-rne.fnl.gov/TBS>

MAL, Pratt, Soltz, Wiedemann
Ann Rev Nucl Part Sci 55 (2005)

mike lisa - Femtoscopy in relativistic heavy ion collisions - Hot Quarks 17 May 2006, Sardinia, Italy

1



Fusion reaction:

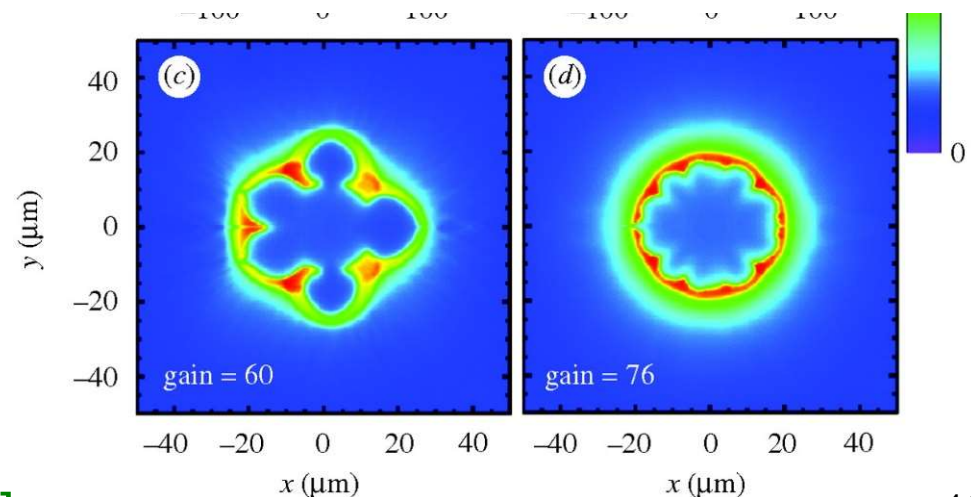


Constant absorptivity,
Spherical irradiation

Ignition temperature = $T_1 \rightarrow$

Simultaneous, volume ignition up to
0.5 R (i.e. **12%** of the volume).

Not too good, but better than:

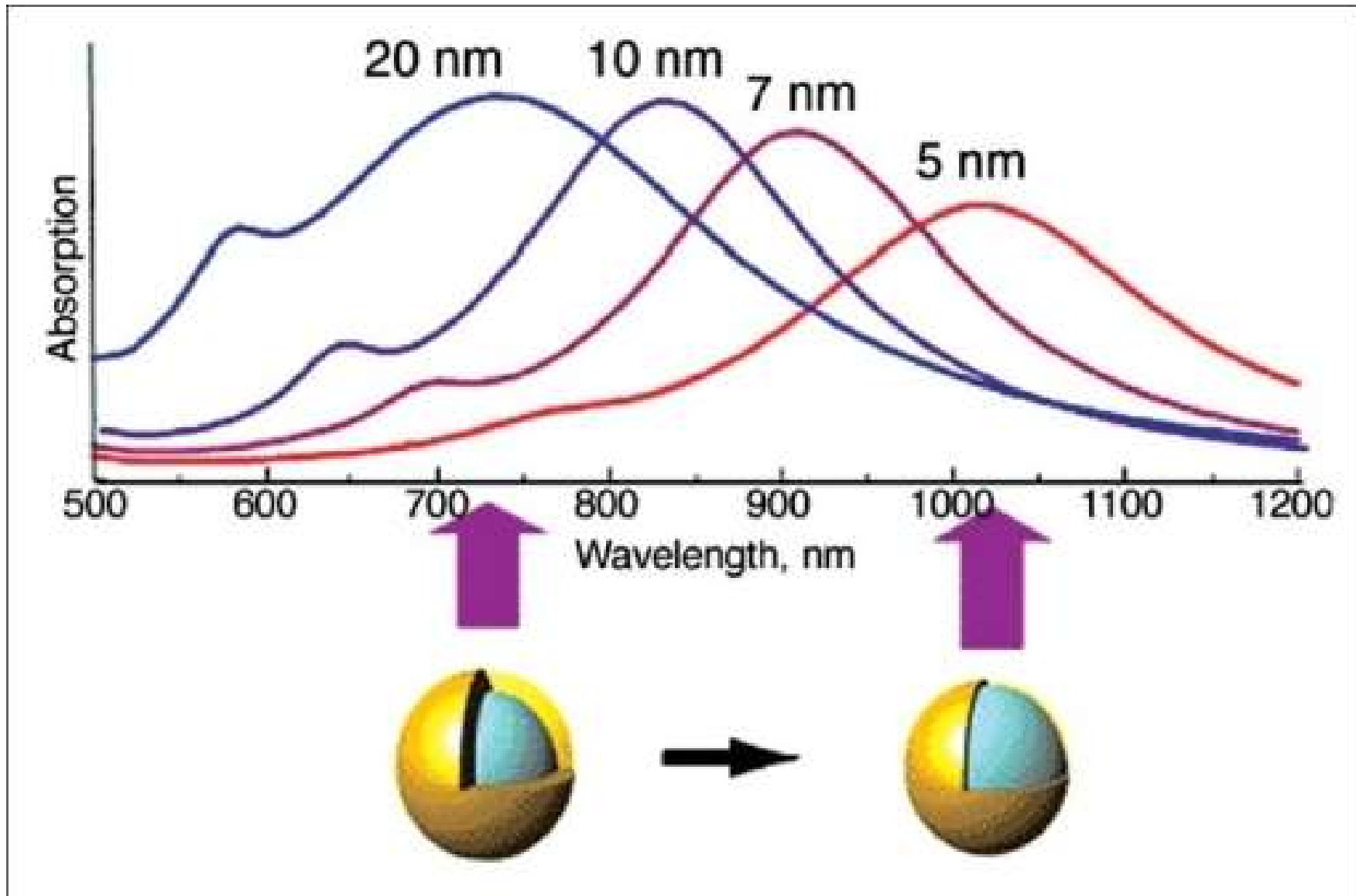


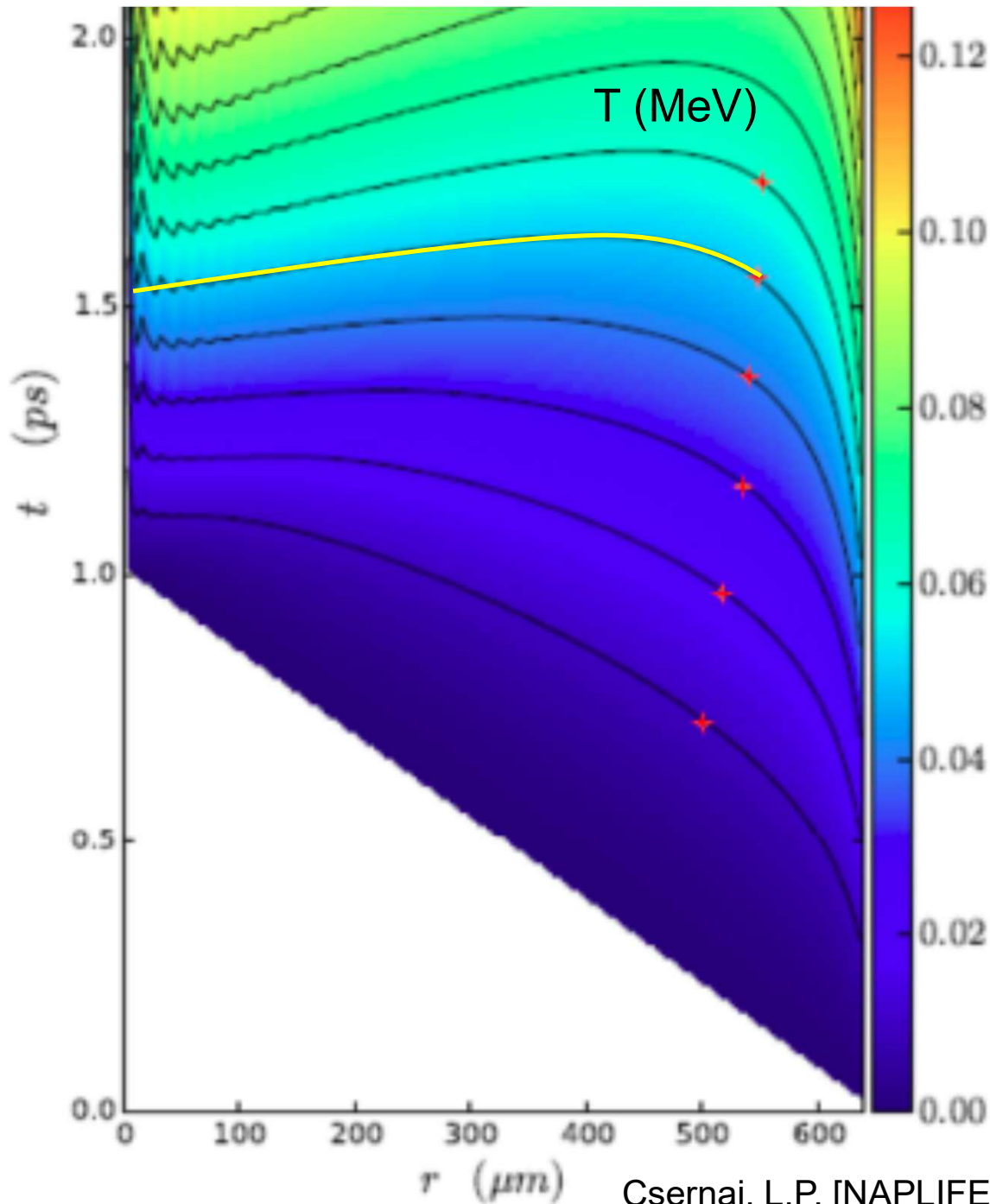
[L.P. Csernai & D.D. Strottman,
Laser and Particle Beams 33, 279 (2015).]

How can we realize it

Idea - #2

Golden Nano-Shells – Resonant Light Absorption

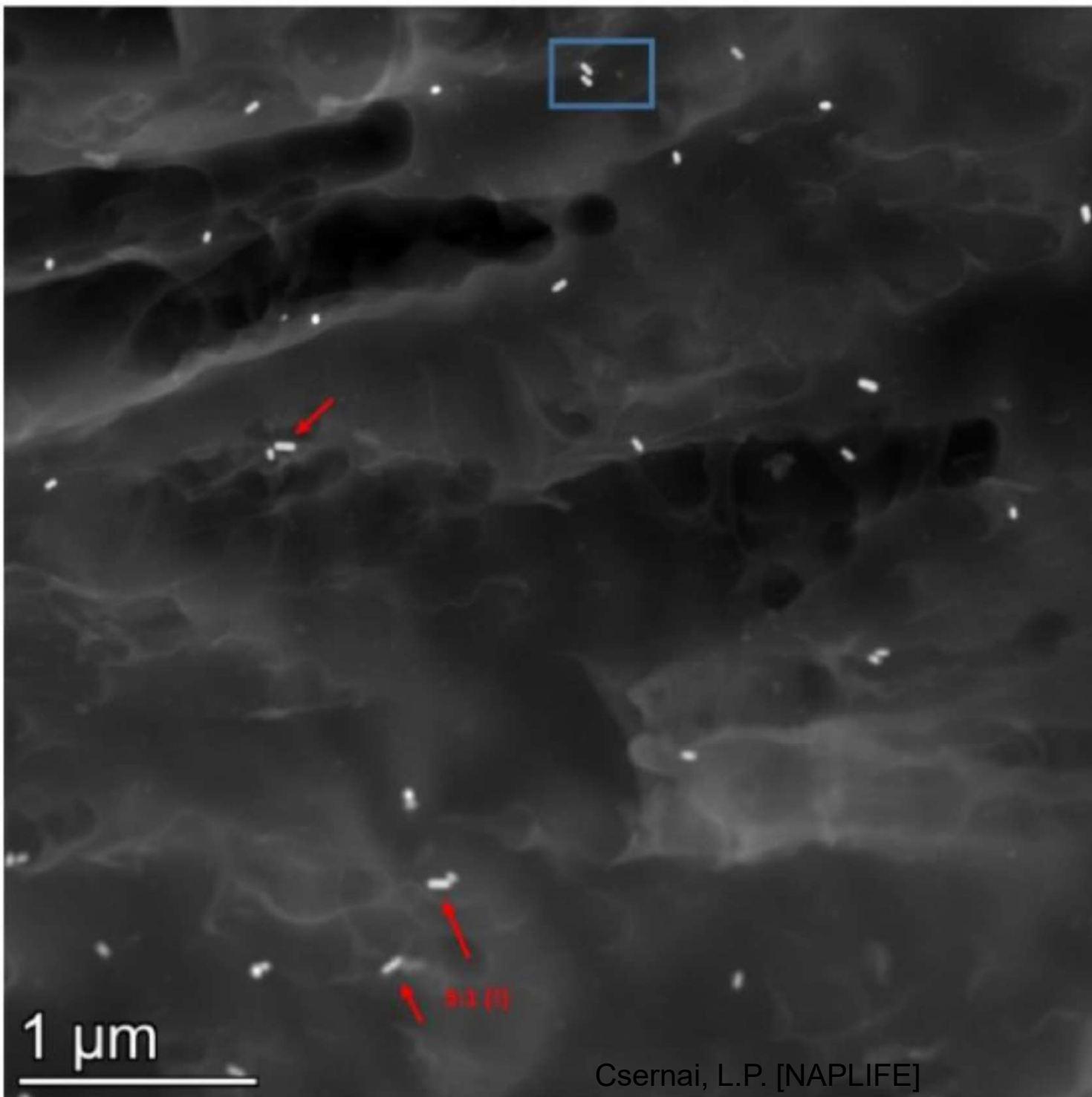




The absorption coefficient is **linearly** changing with the radius: In the center, $r = 0$, $\alpha_K = 30 \text{ cm}^{-1}$ while at the outside edge $\alpha_K = 8 \text{ cm}^{-1}$.

The temperature is measured in units of $T_1 = 272 \text{ keV}$, and $T_n = n T_1$.

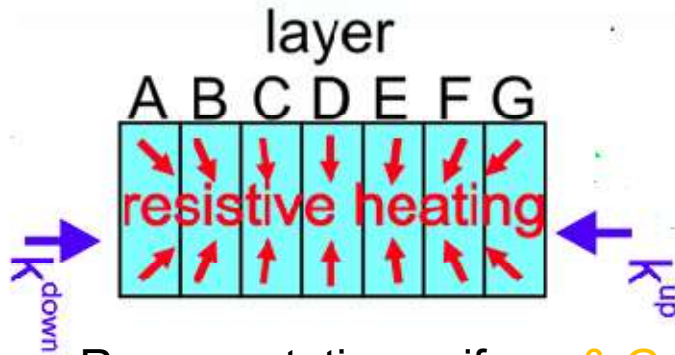
Simultaneous, volume ignition is up to 0.9 R, so 73% of the fuel target!



TEM Photo of
~uniformly
implanted
85x25 nm
nanorod
antennas in
UDMA target
polymer. The
density is
9-20 / μm^3

[**Judit Kámán, A. Bonyár** et al.
(NAPLIFE Collab.),
Gold nanorods ...,
10th ICNFP
2021, **Kolymbari**,
Crete, Greece, 30
August 2021.]¹⁴

Layered target with variable light absorption



Representative uniform & Gaussian number density distributions of (d) 70 oriented nanorods, in a $1 \times 1 \times 21 \mu\text{m}^3$ supercell of UDMA polymer target, with random location distribution.

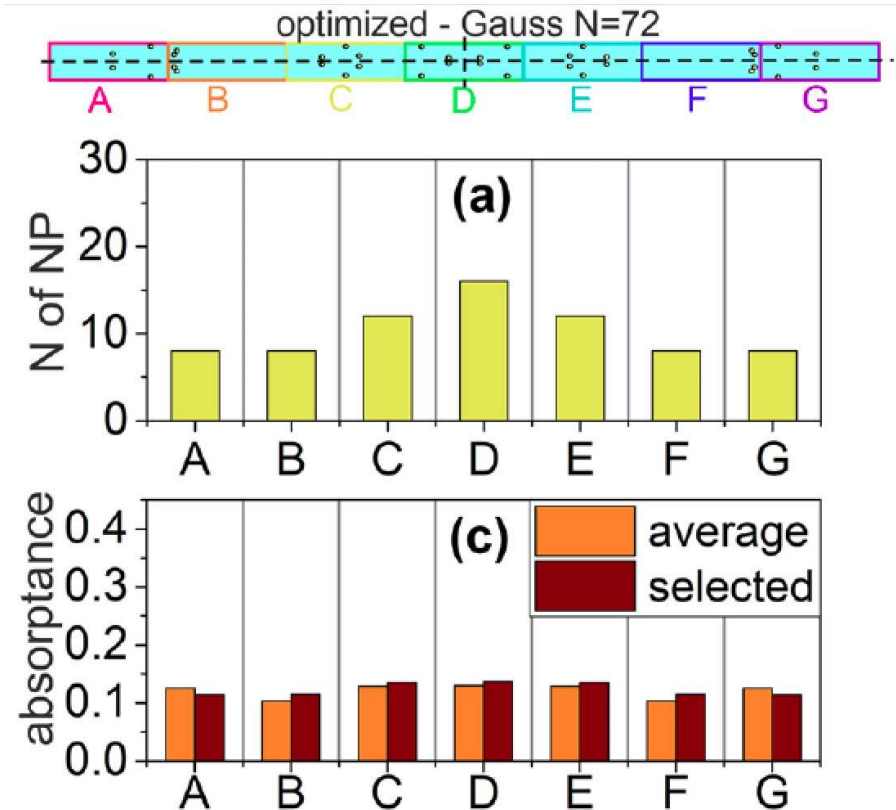
Plasmonics (2022) 17:775–787

<https://doi.org/10.1007/s11468-021-01571-x>

Comparative Study on the Uniform Energy Deposition Achievable via Optimized Plasmonic Nanoresonator Distributions

Mária Csete¹ · András Szenes¹ · Emese Tóth¹ · Dávid Vass¹ · Olivér Fekete¹ · Balázs Bánhelyi² · István Papp^{3,4} · Tamás Bíró³ · László P. Csernai^{3,4,5} · Norbert Kroó^{3,6}

[M. Csete, A. Szenes, E. Tóth, D. Vass, O. Fekete, B. Bánhelyi, T. S. Bíró, L. P. Csernai, N. Kroó: „Comparative study on the uniform energy deposition achievable via optimized plasmonic nanoresonator distributions“, Plasmonics (2022), 17: 775-787; <https://doi.org/10.1007/s11468-021-01571-x>.]



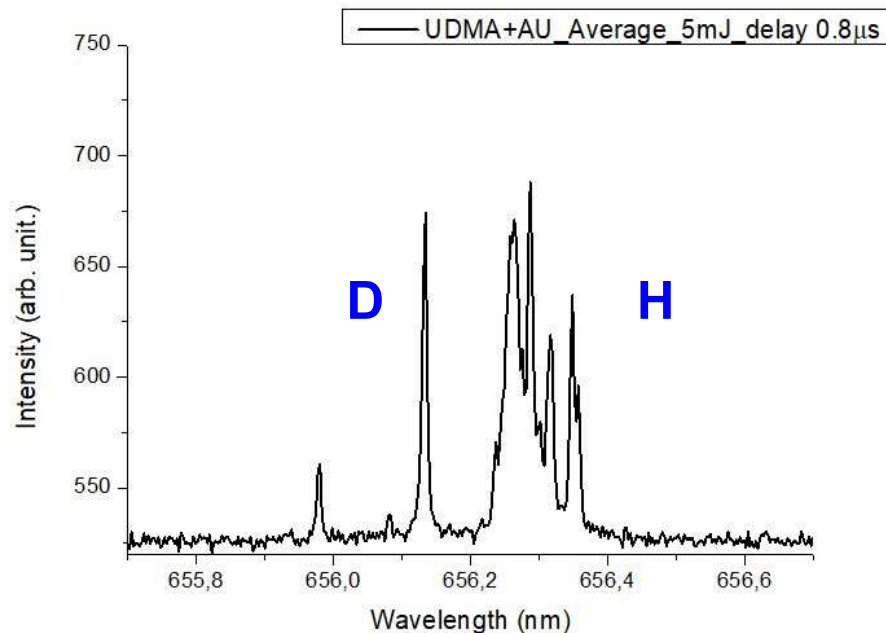
Validation tests →

**Laser Induced Fusion
with Nanoantennas**

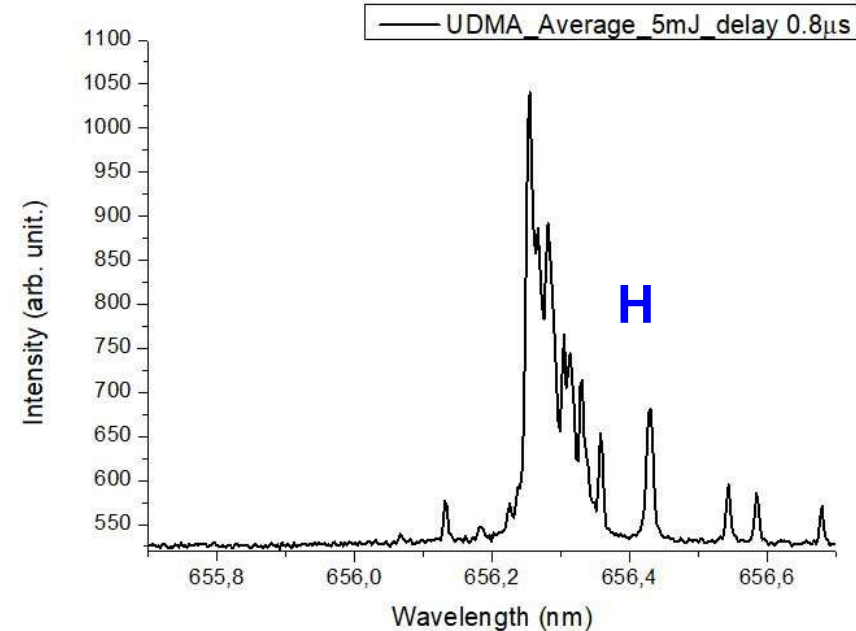
Deuterium production

(PRELIMINARY !)

(N.K.*)



5-12% **D** + 88-95% **H**
~ 10^{17} **D** / pulse (10Hz)



100% **H**
Balmer- α line

Two step process (average of 20 shots):

$p + e^* \rightarrow n + \nu$ \ electron capture (-1.24 MeV)

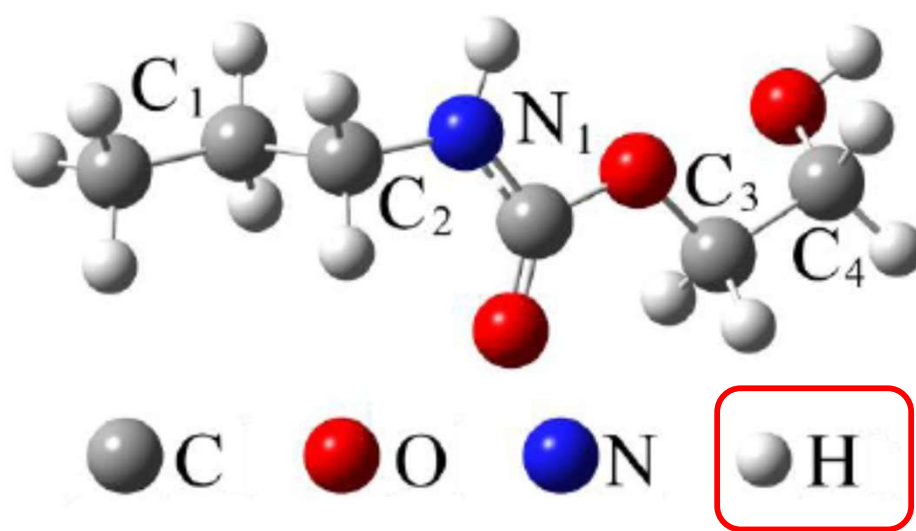
$n + p \rightarrow d + \gamma$ \ neutron capture (+2.22 MeV)

Electron capture may happen spontaneously in heavy nuclei,

here laser light and resonant nanorods act similarly, high e^- density

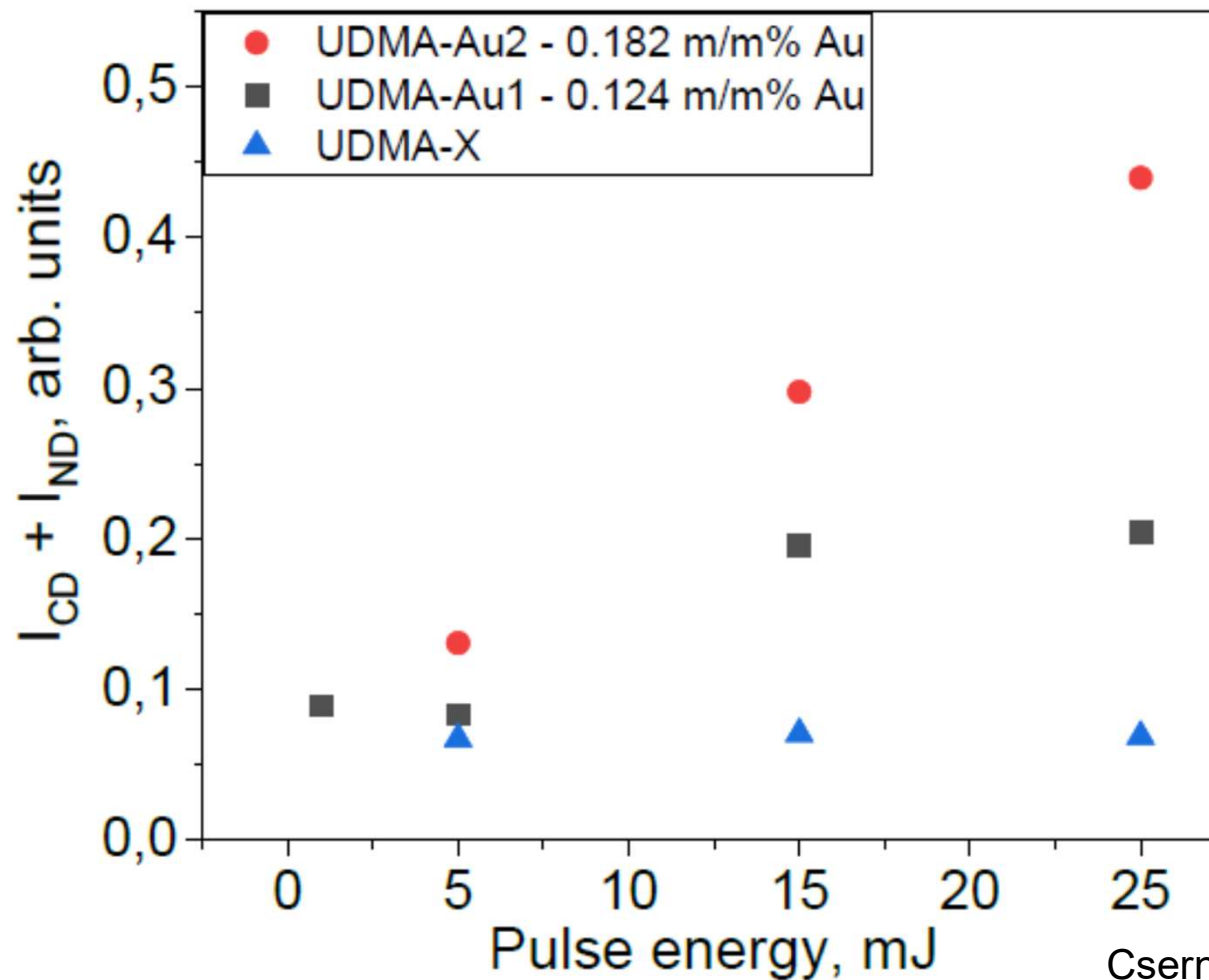
UDMA (470: H38, C23, O8, N2)

[P. Rácz, A. Kumar
et al. 2021 ICNFP,
Kolymbari]



[I. Rigo et al., arXiv 2022]

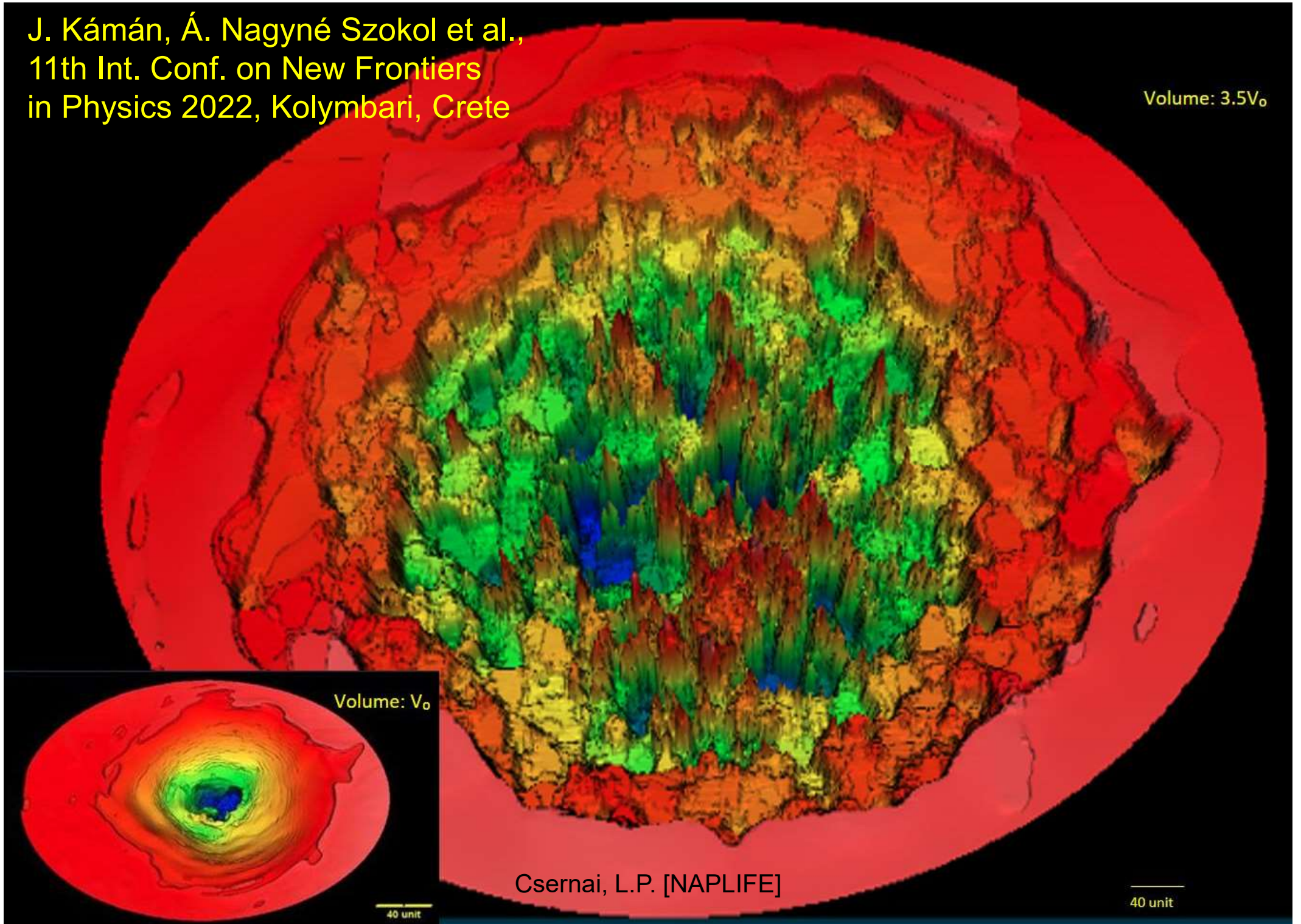
With Nanorods (Au2)
at 25 mJ laser pulse
~4 times increased D
production, compared
to 1 mJ pulse



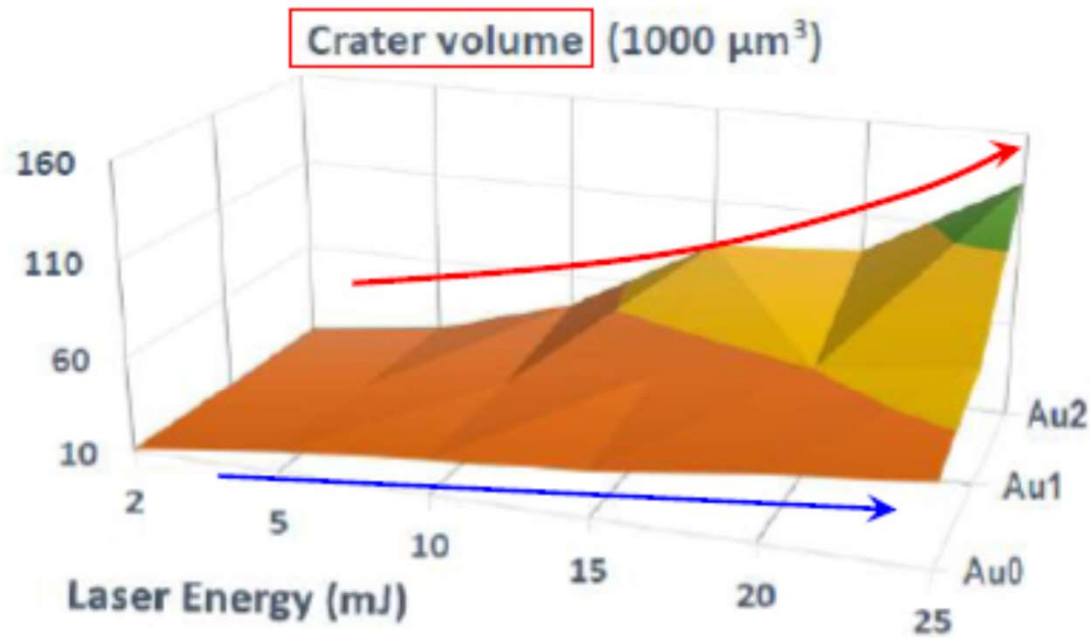
Csernai, L.P. [NAPLIFE]

J. Kámán, Á. Nagyné Szokol et al.,
11th Int. Conf. on New Frontiers
in Physics 2022, Kolymbari, Crete

Volume: $3.5V_0$



Theoretical analysis of Crater & Deuterium production



Crater Formation and Deuterium Production in Laser Irradiation of Polymers with Implanted Nano-antennas

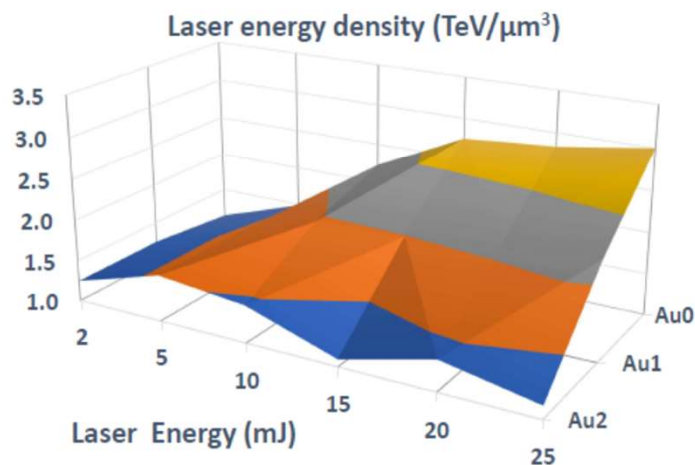
László P. Csernai^{1,2,3}, Igor N. Mishustin³, Leonid M. Sutarov³, Horst Stöcker^{3,7,8}, Larissa Bravina⁴, Mária Coete^{5,6}, Judit Kálmán^{1,5}, Archana Kumari^{1,5}, Anton Motornenko³, István Papp^{1,5}, Péter Rácz^{1,5}, Daniel D. Strottman⁹, András Szenes^{5,6}, Ágnes Szokol^{1,5}, Dávid Vass^{3,6}, Miklós Veres^{1,5}, Tamás S. Biró^{1,5}, Norbert Knö^{1,5,10}
(NAPLIFE Collaboration)

With nanorods V grows non-linearly. Increasing energy deposition. Several types of targets are considered: Au1 and Au2 with implanted nano-rod antennas, and Au0 without implantation. The mass concentrations of implanted particles in UDMA are 0.126% and 0.182% for targets Au1 and Au2, respectively.

With nanorods, Au2, deposited energy into the crater increases non-linearly (!?)

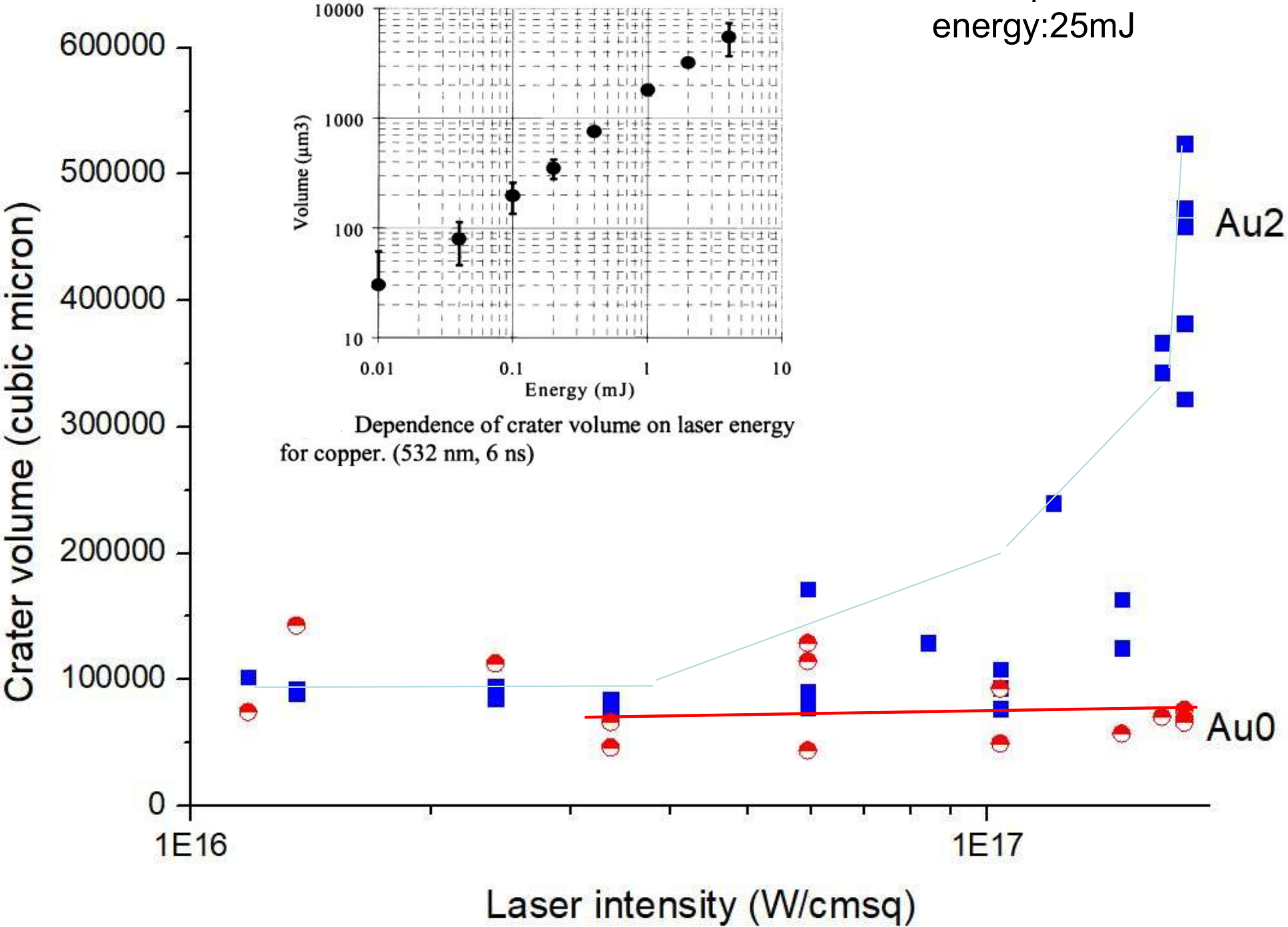
Origin of this extra energy (?)

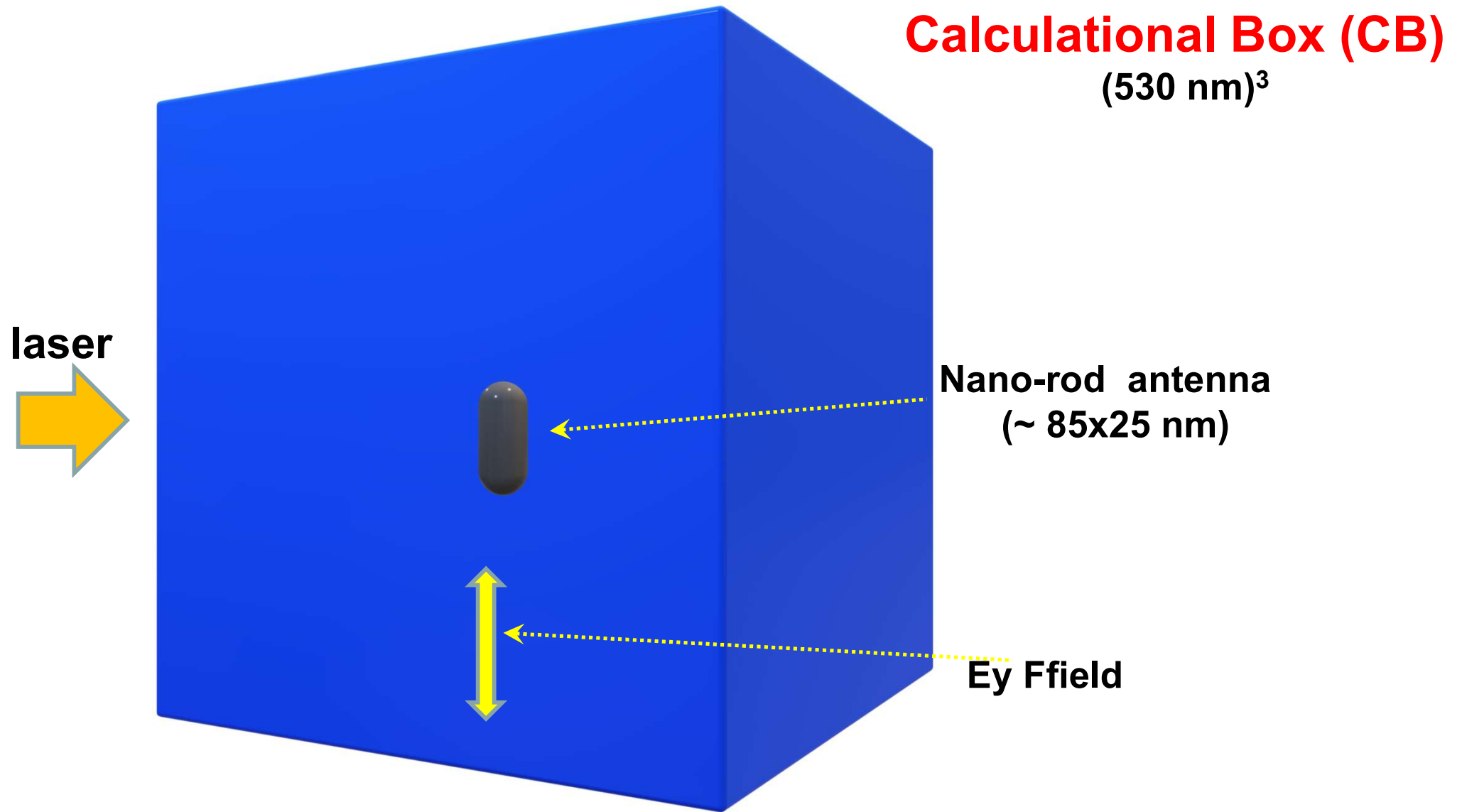
[LP. Csernai et al., Phys. Rev. E, 108(2) 025205 (2023)]



Puzzle ?

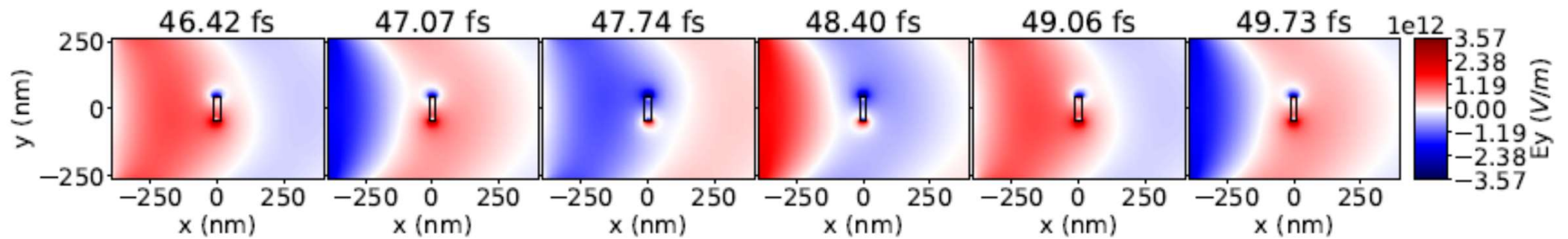
Laser pulse energy: 25mJ



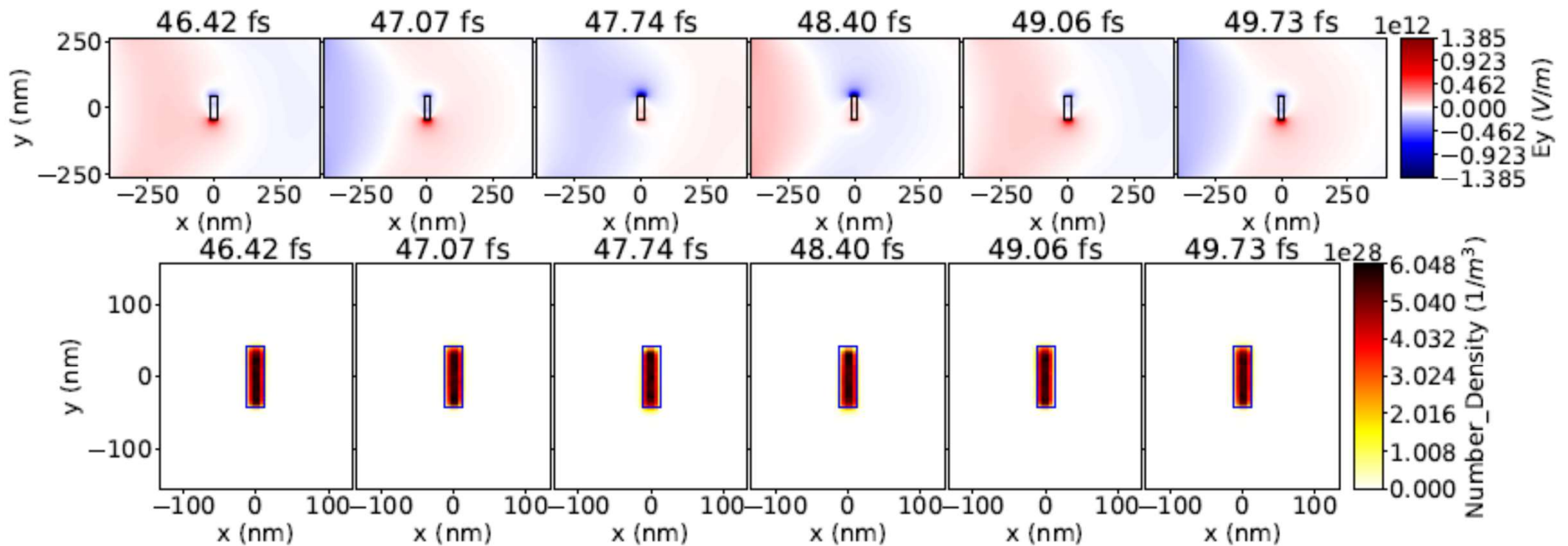


Laser 

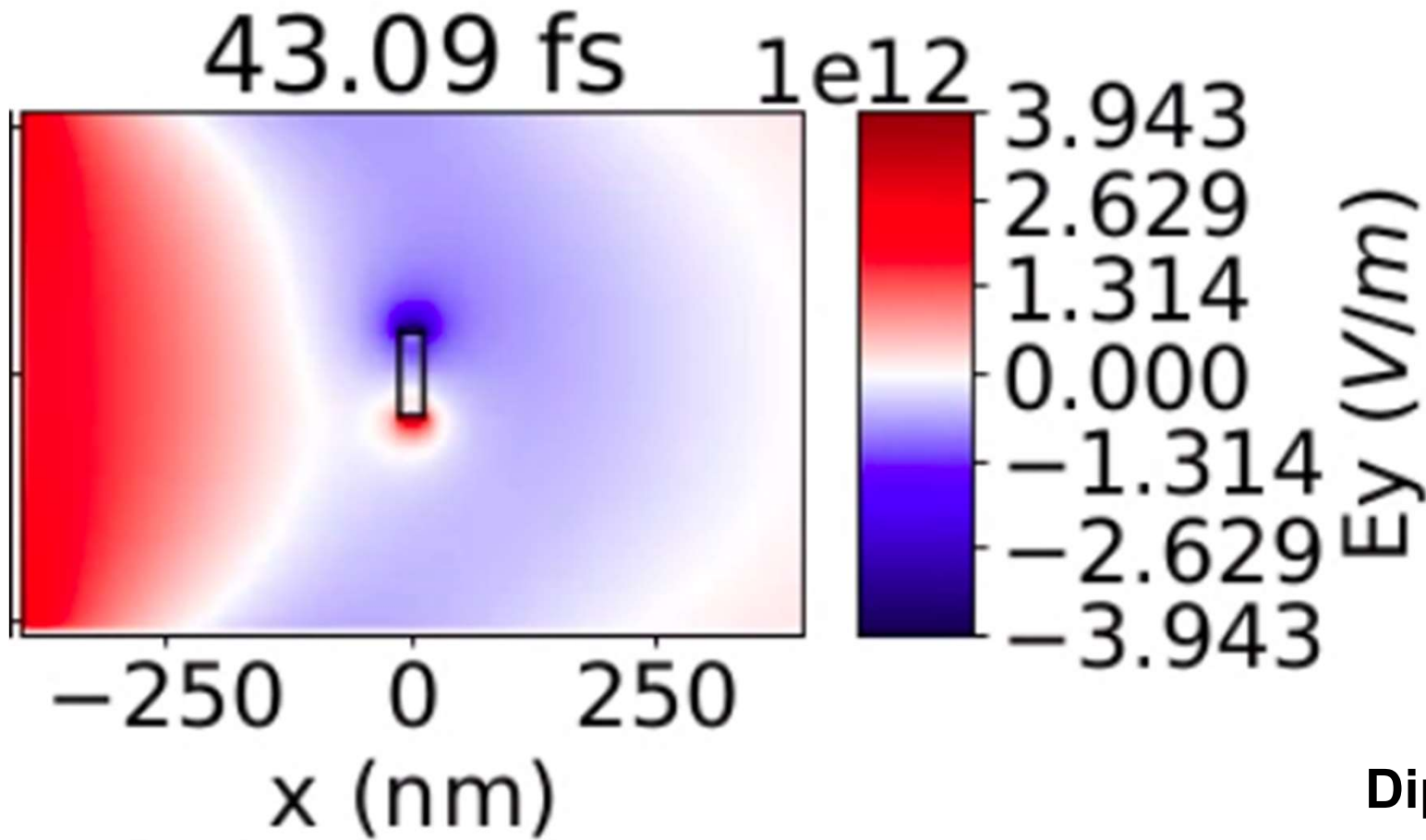
$$I = 4 \cdot 10^{17} \text{ W/cm}^2 \quad V \sim 7.1 \cdot 10^{12} \text{ V/m}$$



$$I = 4 \cdot 10^{15} \text{ W/cm}^2 \quad V \sim 2.6 \cdot 10^{12} \text{ V/m}$$



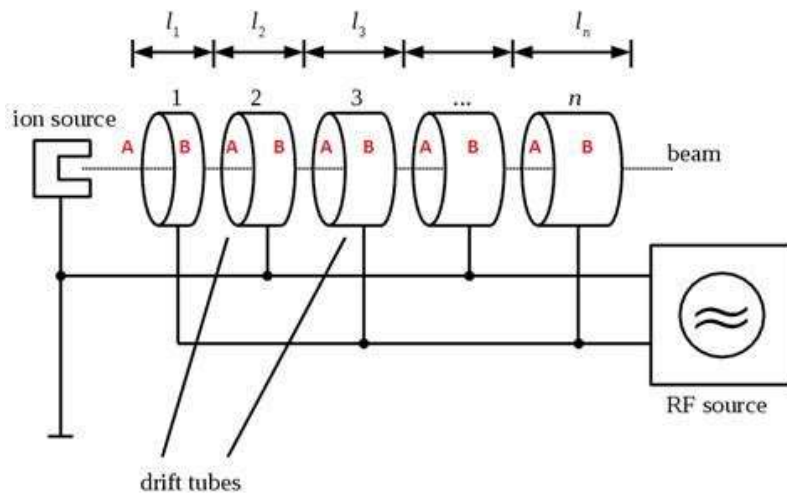
[I. Papp et al., *arXiv*: 2306.13445]



Neighboring protons are accelerated (100-200 nm)

Dipole $L = 85$ nm
 $dV \sim 8 \cdot 10^{12}$ V/m

$$I = 4 \cdot 10^{17} \text{ W/cm}^2$$

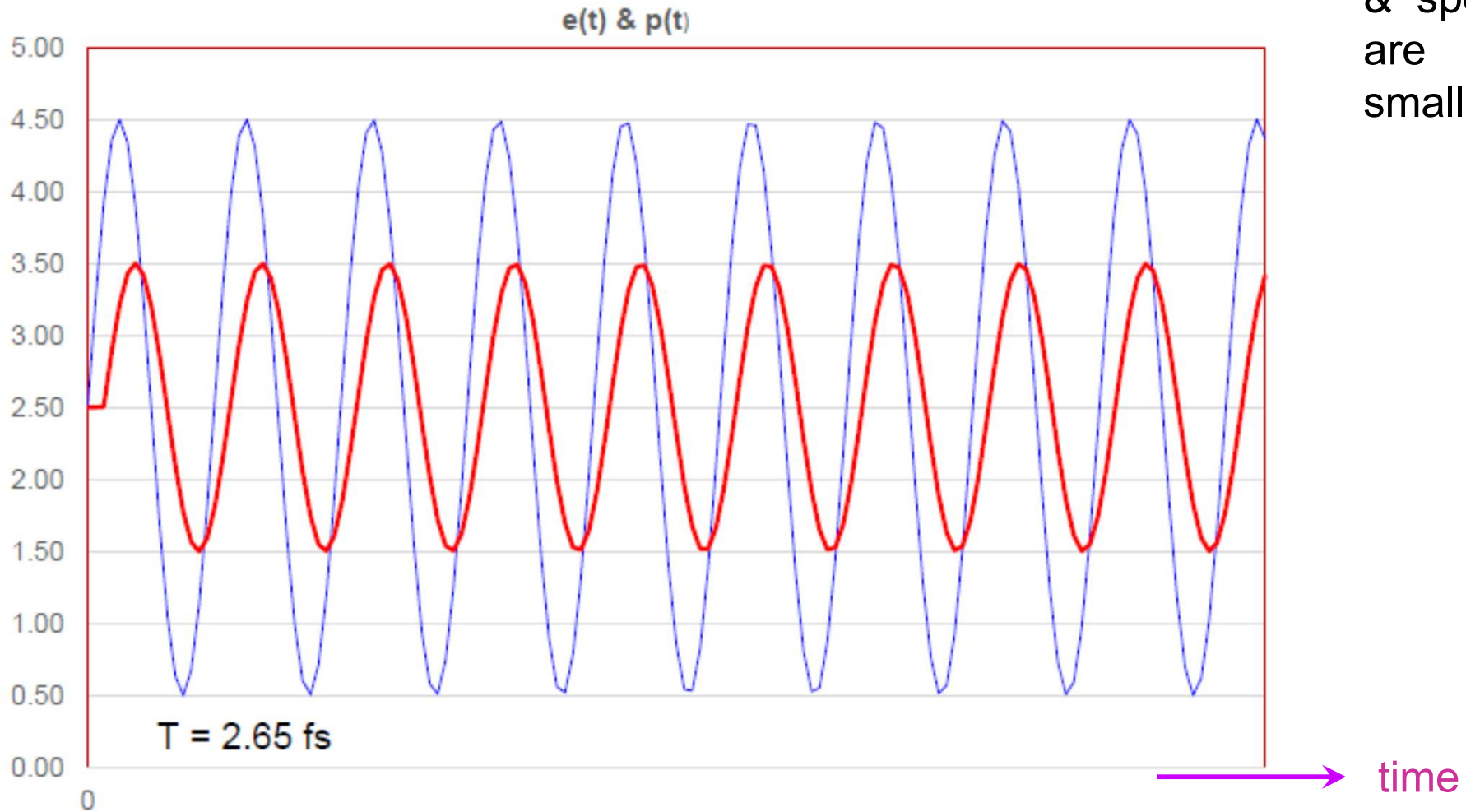


LHC

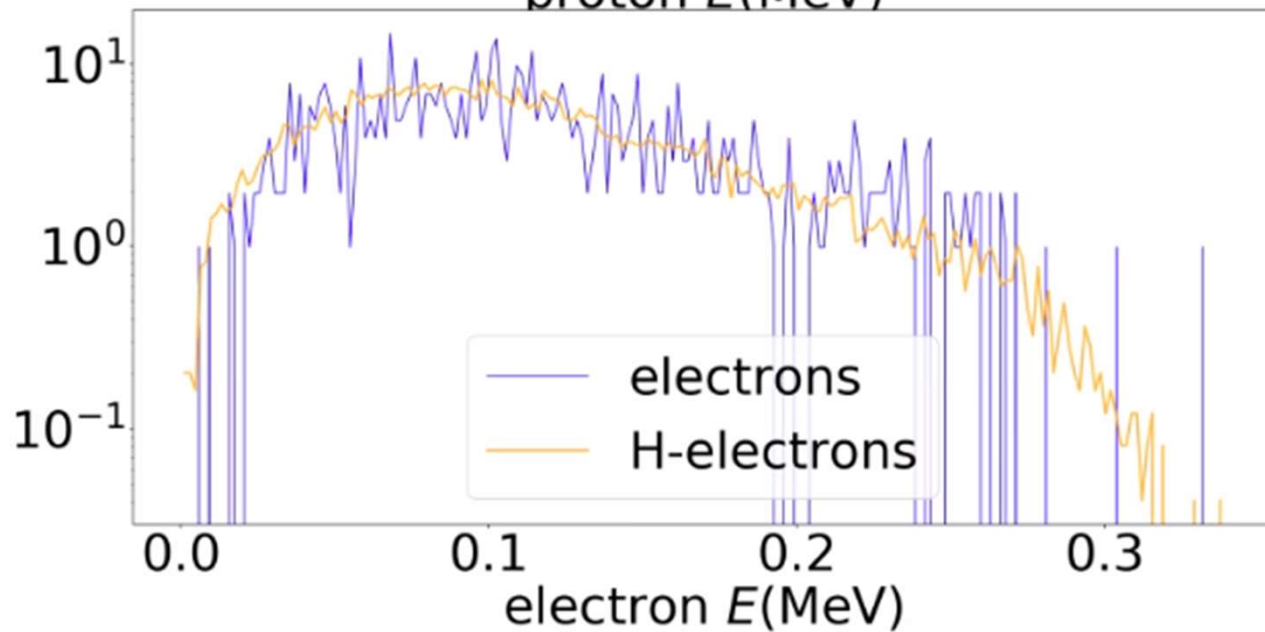
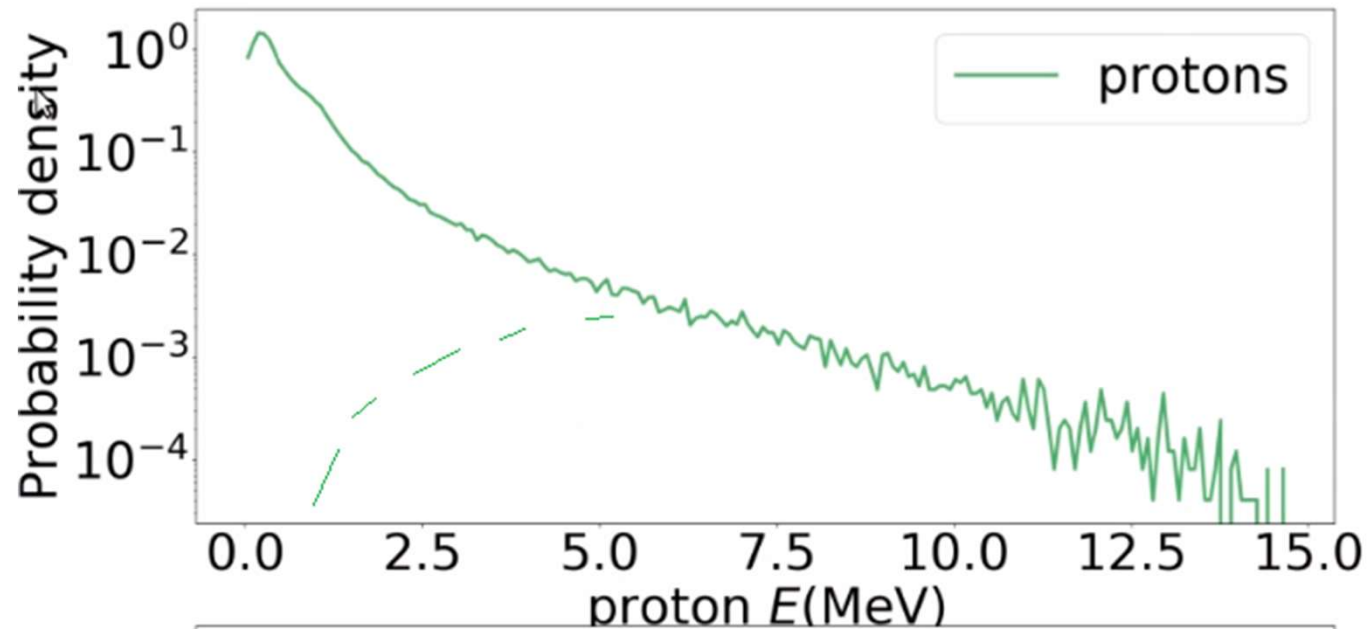
$dV \sim 1 \cdot 10^6$ V/m
 Dipole $L \sim 16$ cm

Laser wake field acceleration mechanism =>

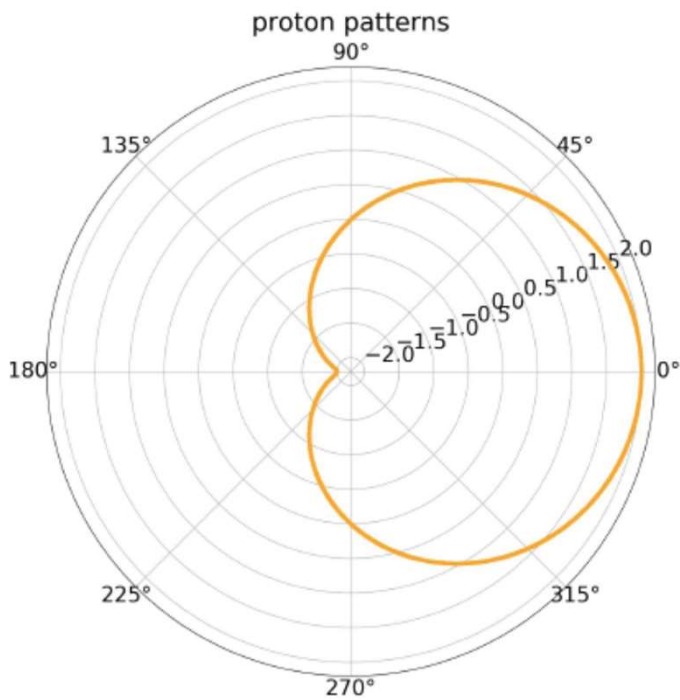
BUT
Proton
amplitudes
& speeds
are
smaller



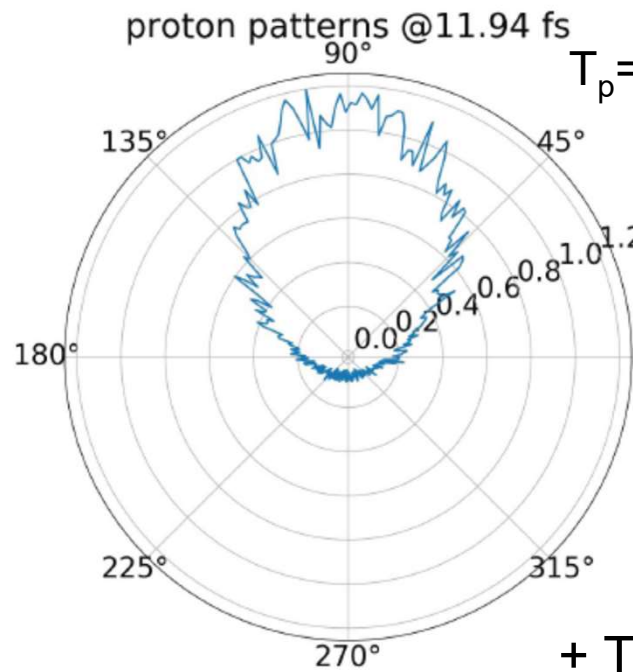
79.56 fs



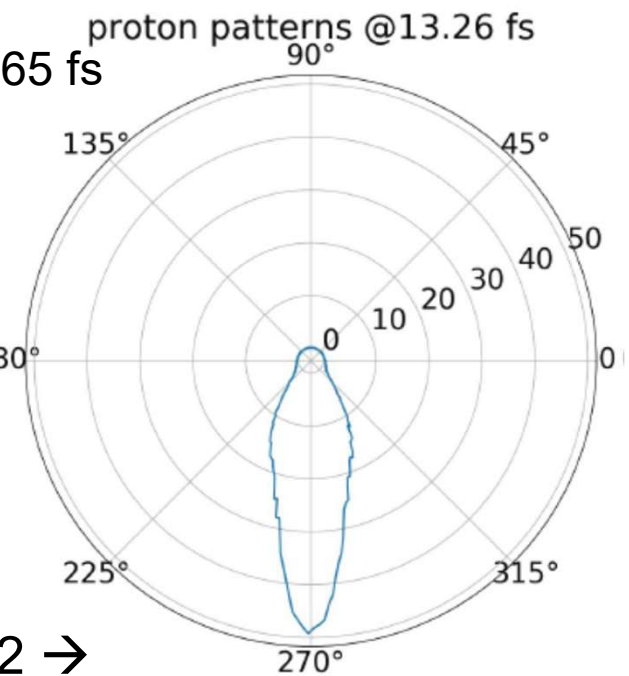
**Number of
1-2 MeV
protons is
about 1-100
=>
small number
of Deuterium**



One-sided irradiation



+ $T_p / 2 \rightarrow$



Two-sided irradiation

$$C(p_1, p_2) = \frac{P_2(p_1, p_2)}{P_1(p_1)P_1(p_2)},$$

$$C(k, q) = 1 + \frac{R(k, q)}{\left| \int d^4x S(x, k) \right|^2},$$

$$R(k, q) = \int d^4x_1 d^4x_2 \cos[q(x_1 - x_2)] \times S(x_1, k + q/2) S(x_2, k - q/2).$$

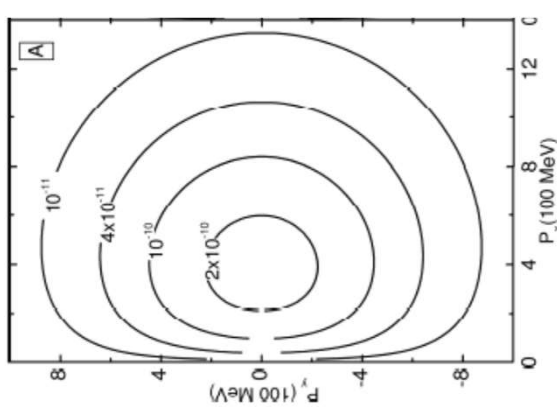
$$J(k, q) = \int d^4x S(x, k + q/2) \exp(iqx)$$

$$R(k, q) = \text{Re} [J(k, q) \overline{J(\vec{k}, -q)}].$$

$$f^J(x, p) = \frac{n(x)}{C_\pi (2\pi\hbar)^3} \exp\left(\frac{-p^\mu u_\mu(x)}{T(x)}\right)$$

$$f_{CJ} = \frac{\Theta(p^\mu d\sigma_\mu) n(x)}{C_\pi (2\pi\hbar)^3} \times$$

$$\left(\exp \frac{-p^\mu u_\mu^R}{T} - \exp \frac{-p^\mu u_\mu^L}{T} \right)$$



Fluid elements [s] can be represented by Cancelling Jüttner distributions, i.e. Cells are not in thermal equilibrium. (protons are accelerated by the nanorods

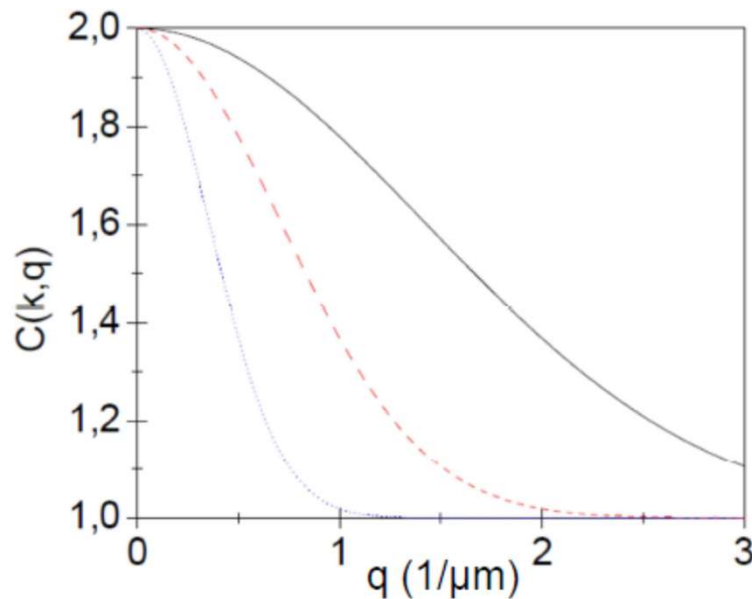
Then standard procedure like for a spherical, thermal cell.

Cancelling Jüttner distribution (CJ)

$$\gamma n_s \exp \left(-\frac{x^2 + y^2 + z^2}{2R^2} \right)$$

$$C(k, q) = 1 + \exp \left(-(\Delta\tau)^2 (\hat{\sigma}^\mu q_\mu)^2 - R^2 q^2 \right)$$

$$C(k, q) = 1 + \exp \left(-R^2 q^2 \right)$$



For a single spherical source:

[L.P. Csernai, S. Velle, and D.J. Wang
PHYSICAL REVIEW C 89, 034916 (2014)]

FIG. 6. (color online) The correlation function, $C(k, q)$, for a single, static, spherically symmetric, Gaussian source with different radii, $R = 4, 1$ and $0.25 \mu\text{m}$, (*blue dotted, red dashed, and full black lines* respectively), as described by Eq. (16).

New fusion mechanism

Traditionally (NIF) after ignition, DT burning is spreading by *alpha particle self heating*. This turns out to be slower than expansion after extreme compression and extreme pressure.

HINT:

Here after simultaneous (time-like) ignition attraction of large number of electrons *collectively accelerate* protons, which can induce nuclear reaction (e.g. transmutation).

We try to verify this mechanism by the Hanbury-Brown and Twiss effect to determine the deuteron and alpha source size, after a laser shot.

**High Energy, Short Pulse Laser,
unique
at
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European Laser Infrastructure – Szeged, HU



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EU Extr. Light Infrastructure
Attosec. Light Pulse Source

2PW High Field laser
10 Hz, <10fs, **20 J**