

Relativisztikus Lézer-fúzió Nanotechnológiával (Relativistic Laser Fusion with Nanotechnology)



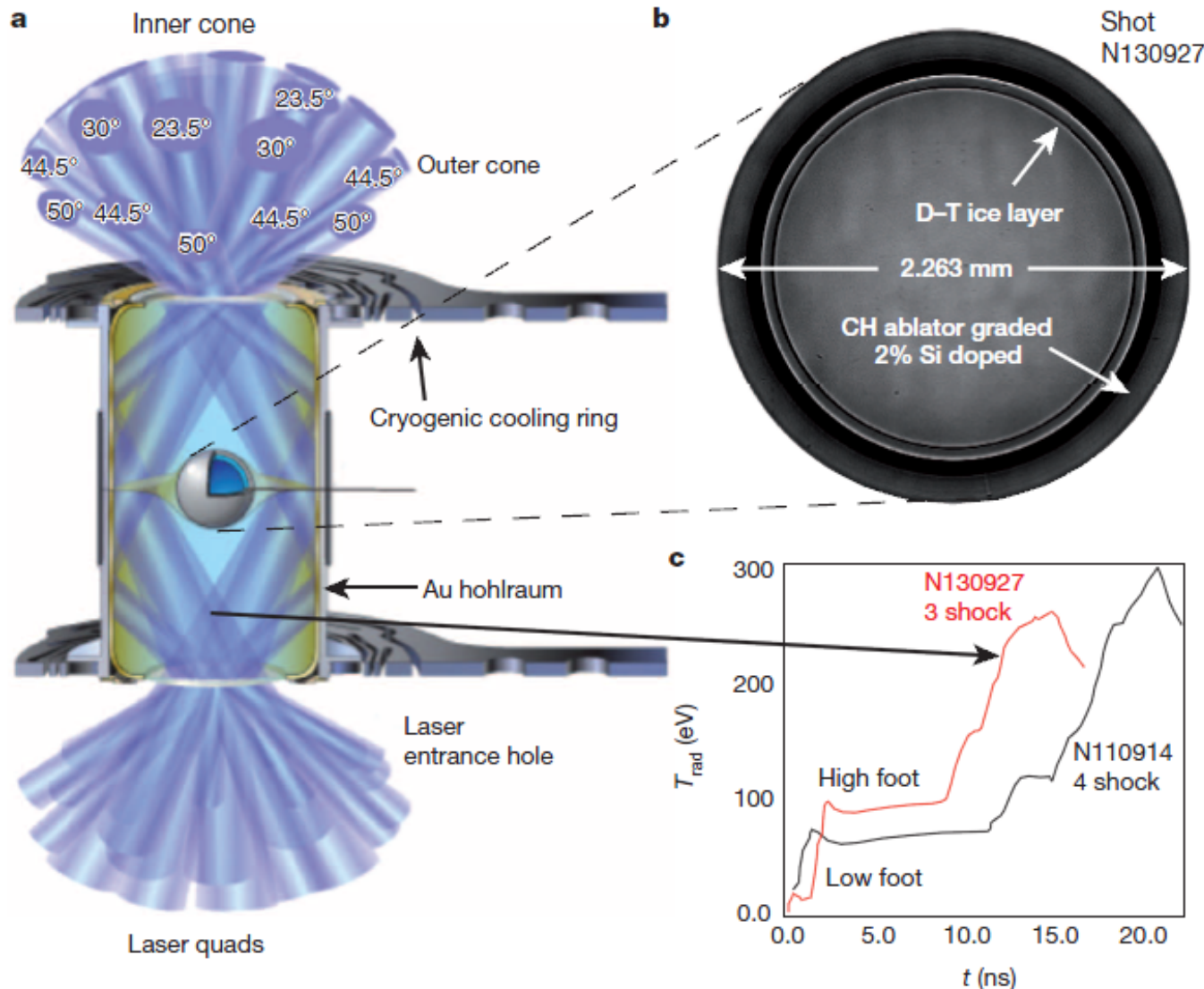
1973

Csernai László, Bergeni Egyetem

Fúziós forum, February 11, 2020 , Budapest

Dept. of State Secretary for the Capacity Maintaining of Paks Nuclear power Station

Indirectly Driven, ICF target for NIF



Time profile of the laser beam:

Initial pre-compression of \sim **10 ns**,
→ Stable compression

→ Then final “shocks” of \sim **15 ns** to ignite

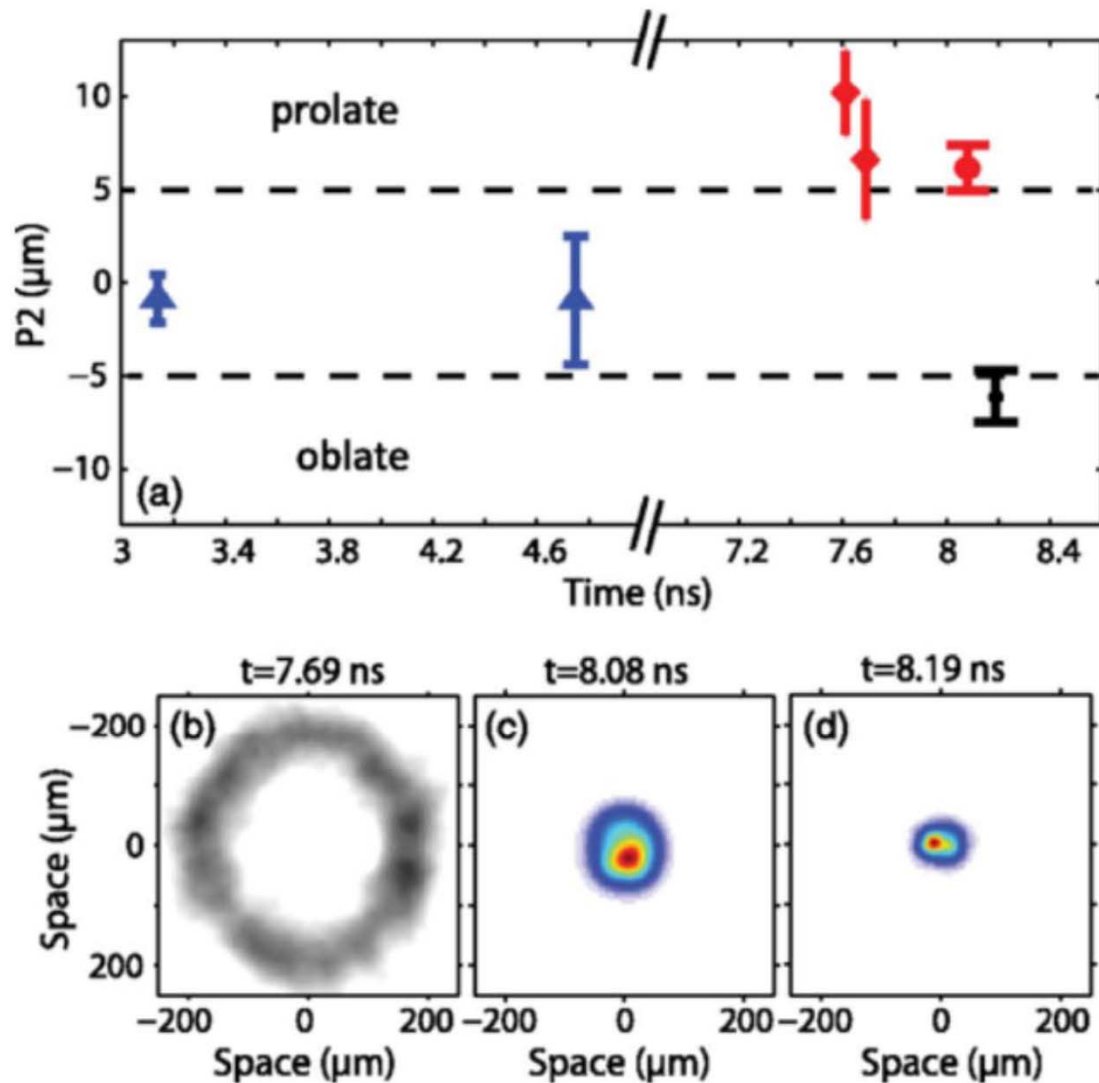


Figure 2

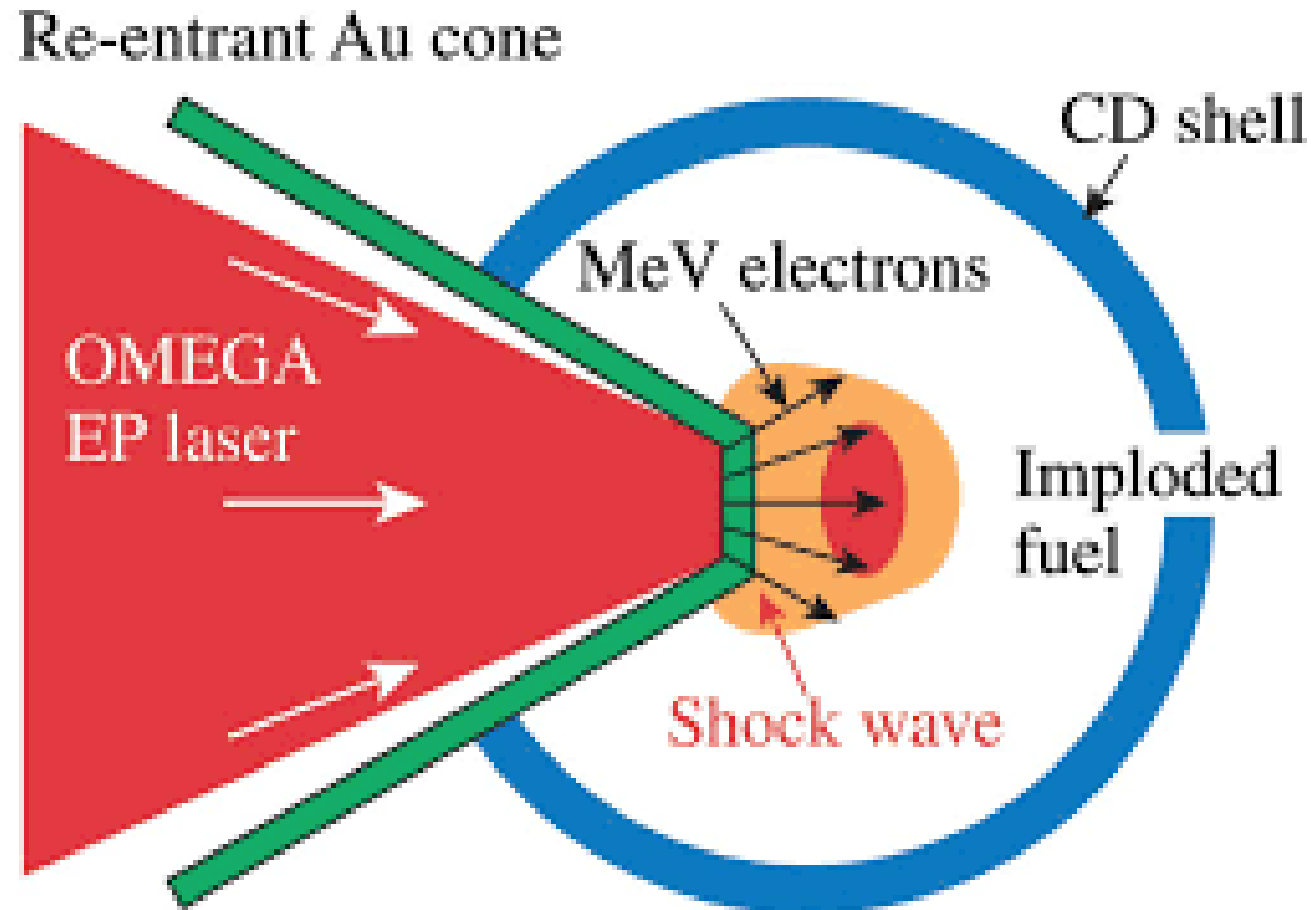
(a) History of the implosion symmetry for the doped HDC capsule measured at increasing convergence and time using a succession of experimental techniques. Blue points are keyhole data, red points are 2DconA data, black point is the DT cryogenic platform. The definition of $P2$ in microns as a measure of deviation from round is described in the text. (b) equatorial x-ray radiograph of the shell, (c) equatorial x-ray image of the hot spot at bang time (convergence 17) (d), equatorial x-ray image of the hot spot at bang time (convergence 25).

Notice: The ignition peak is now in the centre of the compressed target pellet!

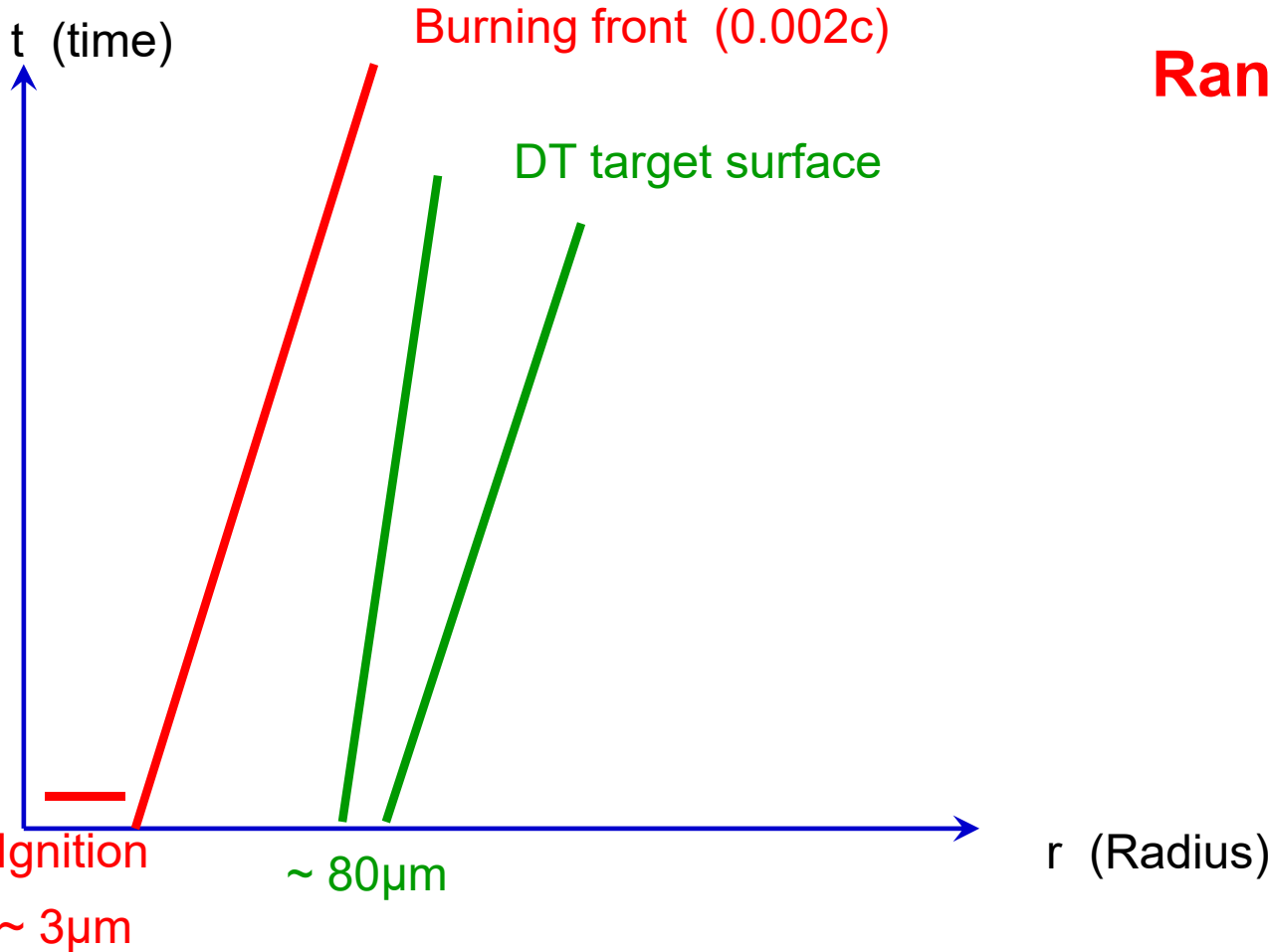
80 μm - 2018

OMEGA Laser Facility at the University of Rochester in New York

80 Lasers

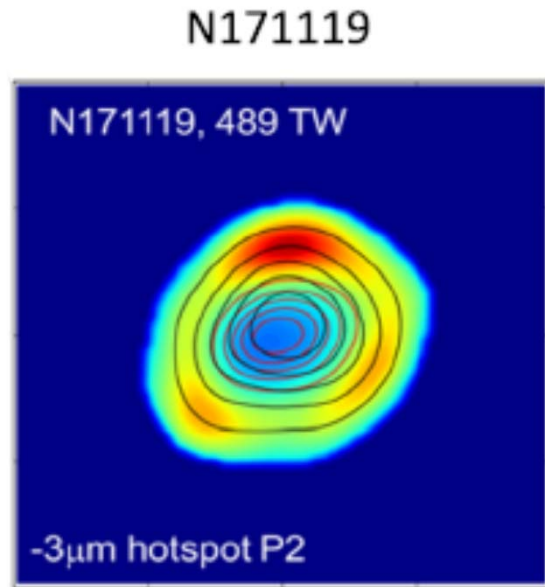


Classical Fluid Dynamics – Burning front speed



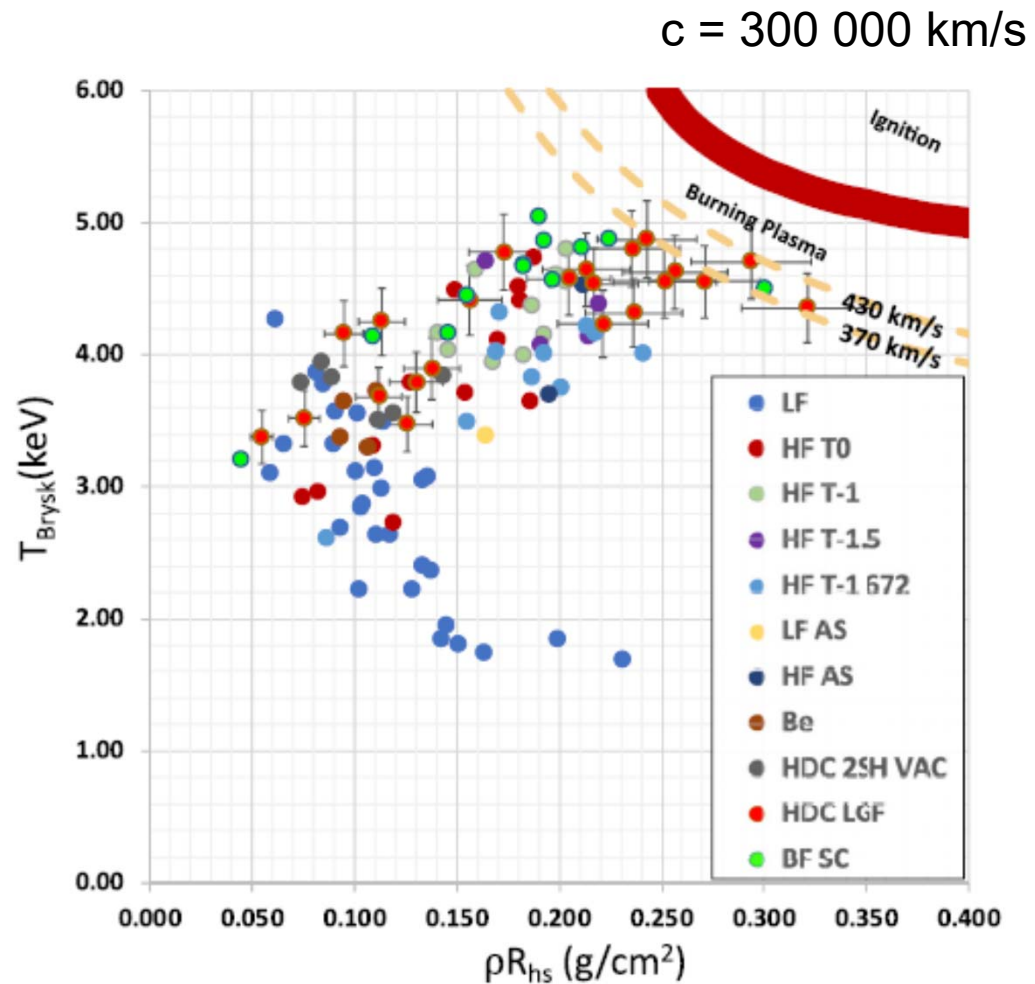
**Rankine-Hugoniot
relation**

LLNL – NIF April 2019



Hot-spot, 3 μ m, burning spreads by α -heating: boundary is approached.

[O.A. Hurricane et al., Phys. Plasmas 26, 052704 (2019) April.]



[A.H. Taub (1948)]

PHYSICAL REVIEW

VOLUME 74, NUMBER 3

AUGUST 1, 1948

Relativistic Rankine-Hugoniot Equations

A. H. TAUB

*University of Illinois, Urbana, Illinois and Institute for Advanced Study, Princeton University, Princeton, New Jersey**

Next we suppose that the three-dimensional volume is a shell of thickness ϵ enclosing a surface of discontinuity Σ whose three-dimensional normal vector is Λ_i . If we choose our coordinate system so that the discontinuity is at rest, then since

$$\underline{\lambda_\alpha \lambda^\alpha = 1}, \quad \sum_{i=1}^3 \Lambda_i^2 = 1,$$

we have

$$\lambda_i = \Lambda_i \quad \text{and} \quad \underline{\lambda_4 = 0}.$$

Hence Eqs. (7.1) and (7.2) become, as ϵ goes to zero,

$$[\rho^0 u^i \Lambda_i] = 0, \quad (7.3)$$

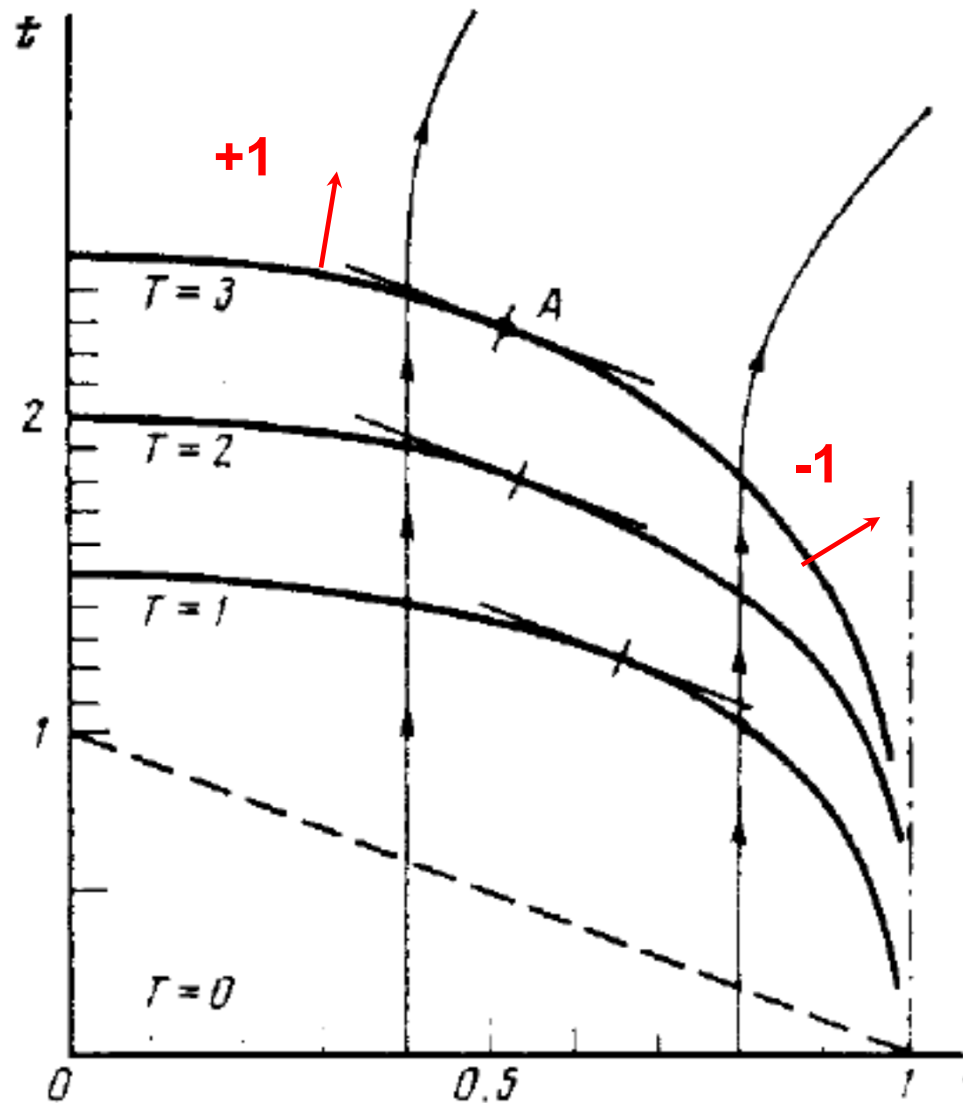
$$[T^{\alpha i} \Lambda_i] = 0, \quad (7.4)$$

where

$$[f] = f_+ - f_-$$

Taub assumed that (physically) only slow space-like shocks or discontinuities may occur (with space-like normal, $\lambda_4=0$).

This was then taken as standard, since then (e.g. LL 1954-)



[L. P. Csernai, Zh. Eksp. Teor. Fiz. 92, 379-386 (1987) & Sov. Phys. JETP 65, 216-220 (1987)]

corrected the work of [A. Taub, Phys. Rev. 74, 328 (1948)]

$$\lambda_\alpha \lambda^\alpha = \pm 1$$

Л. П. Чернаи

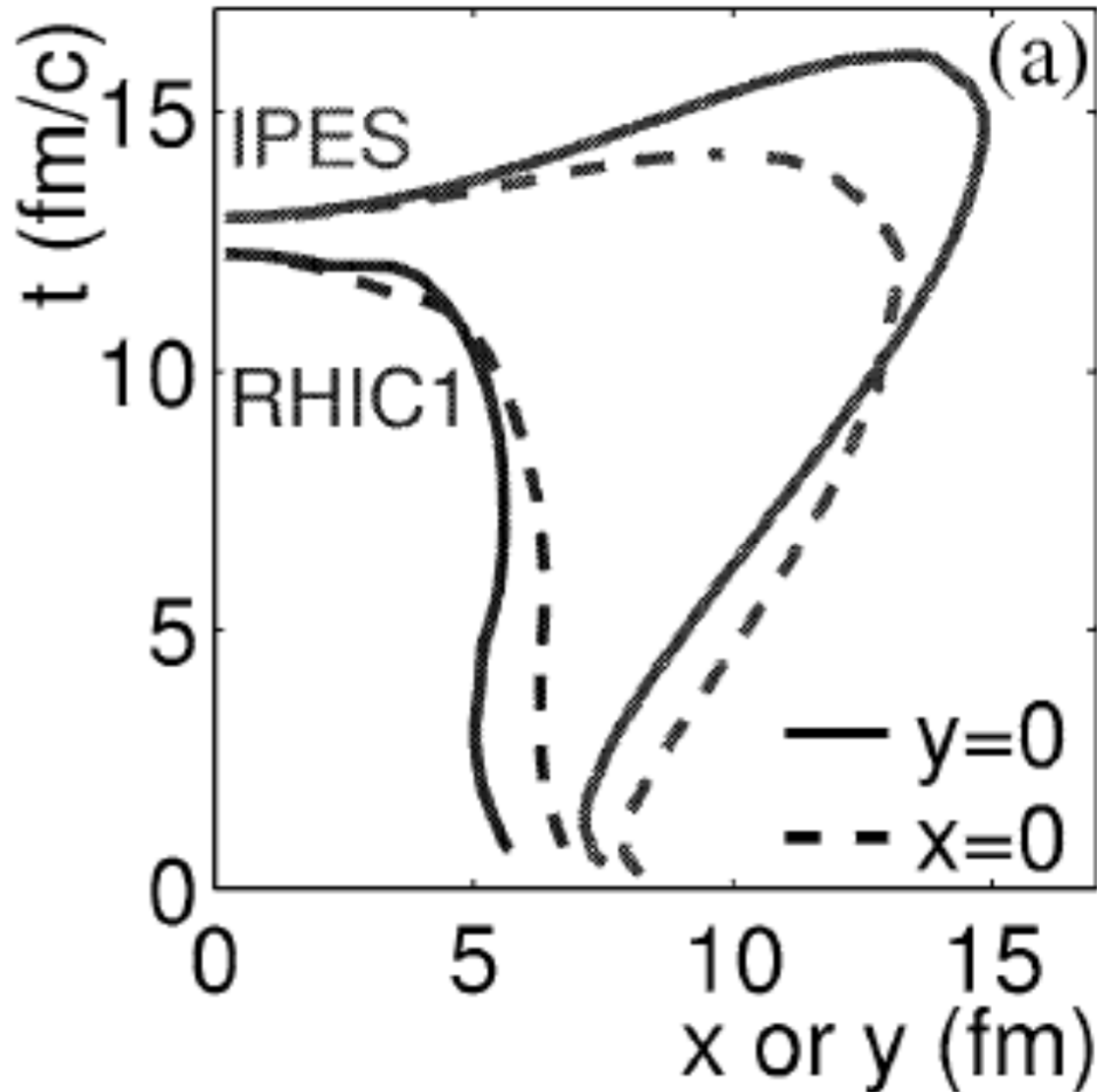
ДЕТОНАЦИЯ НА ВРЕМЕНИПОДОБНОМ ФРОНТЕ
ДЛЯ РЕЛЯТИВИСТСКИХ СИСТЕМ

Журнал экспериментальной и теоретической физики

8

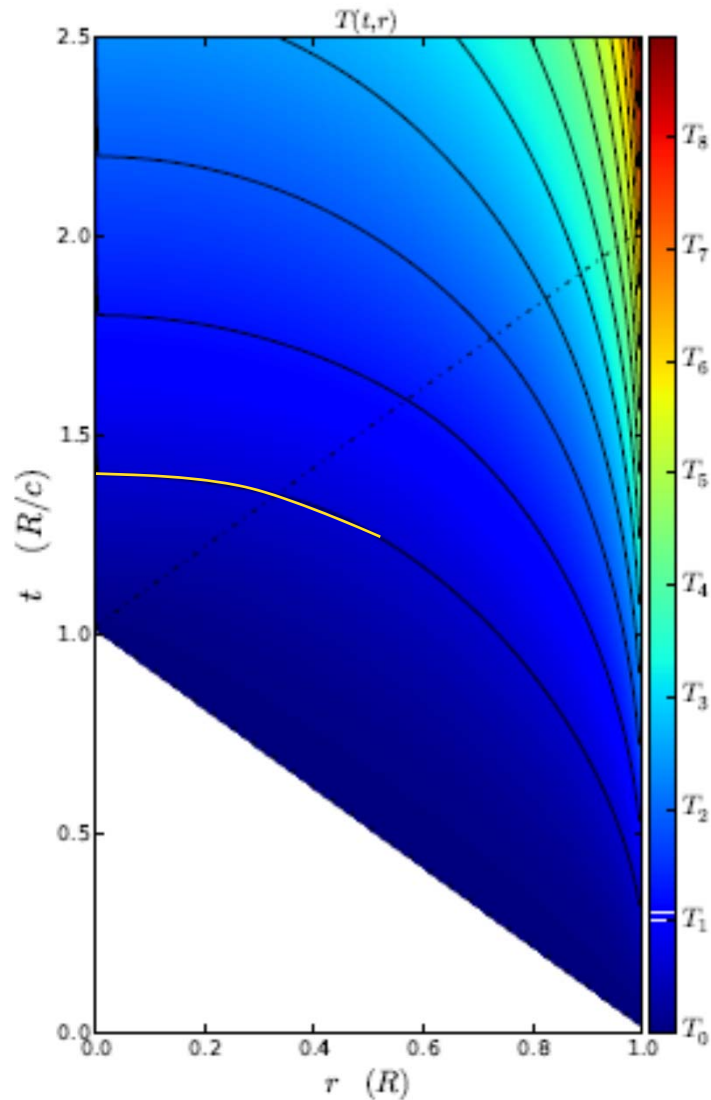
1987

Relativistic burning of Quark-Gluon Plasma



Discovery of QGP:
2000 CERN
2001 BNL

[U.W. Heinz and
P.F. Kolb, Phys.
Lett. B 542, 216
(2002)]



Fusion reaction:



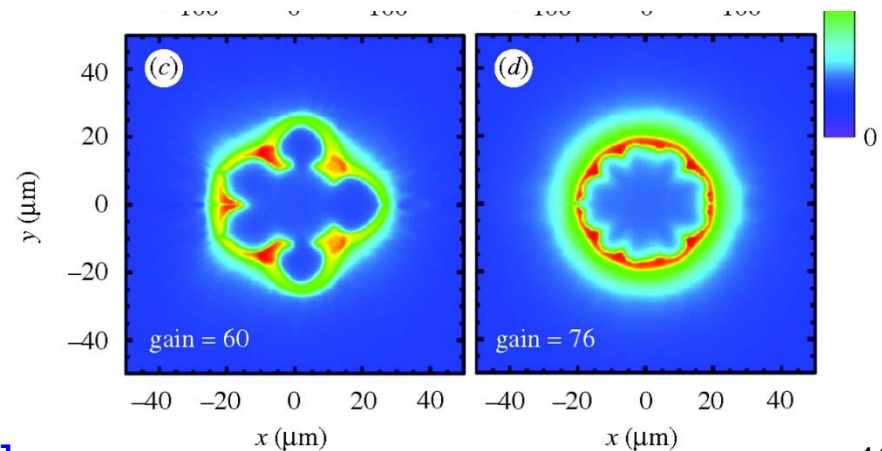
Constant absorptivity,

Spherical irradiation

Ignition temperature = $T_1 \rightarrow$

Simultaneous, volume ignition up to
 $0.5 R$ (i.e. **12%** of the volume).

Not too good, but better than:



[L.P. Csernai & D.D. Strottman,
 Laser and Particle Beams 33, 279 (2015).]

Can we achieve larger volume ignition (II-nd)

Two ideas are combined by

L.P. Csernai, N. Kroo, I. Papp [**Patent # P1700278/3**](*)

- Heat the system uniformly by radiation with **RFD**
- Achieve uniform heating by **Nano-Technology**

Mechanical compression and adiabatic heating should be reduced, because it is slow and leads to Rayleigh-Taylor instabilities. Similarly outside ablator surface should be reduced also.

Uniform, **4π radiation** should heat the target to ignition within the light penetration time (i.e. ~ 10 - 20 ps). This follows from RFD!

[L.P. Csernai, N. Kroo, I. Papp, *Laser and Particle Beams*,
LPB, 36(2), (2018) 171-178. .
<https://doi.org/10.1017/S0263034618000149>]

LPB, 36(2), (2018) 171-178.

Laser and Particle Beams

cambridge.org/lpb

Research Article

Cite this article: Csernai LP, Kroo N, Papp I (2018). Radiation dominated implosion with nano-plasmonics. *Laser and Particle Beams* 1–8. <https://doi.org/10.1017/S0263034618000149>

Received: 28 November 2017

Revised: 14 March 2018

Accepted: 3 April 2018

Key words:

Inertial confinement fusion; nano-shells; relativistic fluid dynamics; time-like detonation

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... and 35th Hirschegg
Int. Workshop on High
Energy Density
Physics, Jan. 25-30,
2015

Radiation dominated implosion with nano-plasmonics

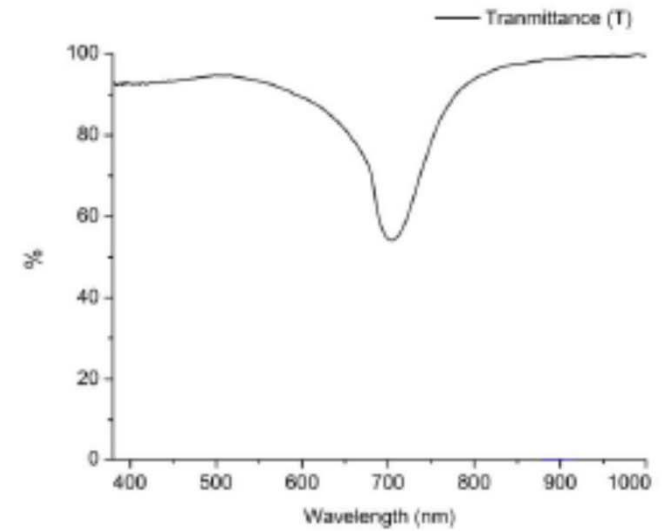
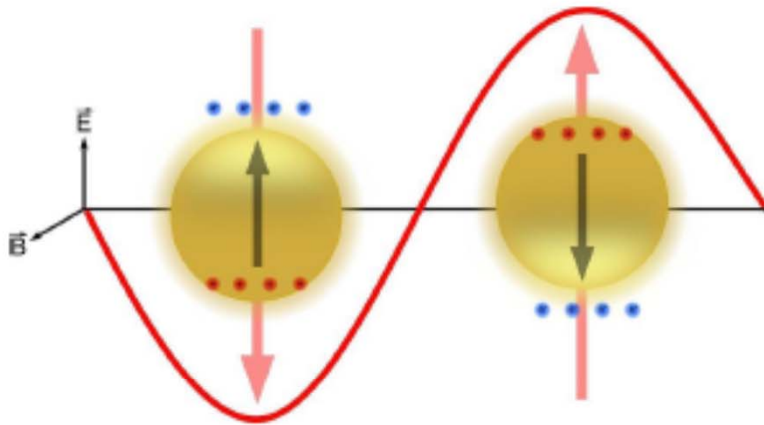
L.P. Csernai¹, N. Kroo^{2,3} and I. Papp⁴

¹Department of Physics and Technology, University of Bergen, Bergen, Norway; ²Hungarian Academy of Sciences, Budapest, Hungary; ³Wigner Research Centre for Physics, Budapest, Hungary and ⁴Department of Physics, Babes-Bolyai University, Cluj, Romania

Abstract

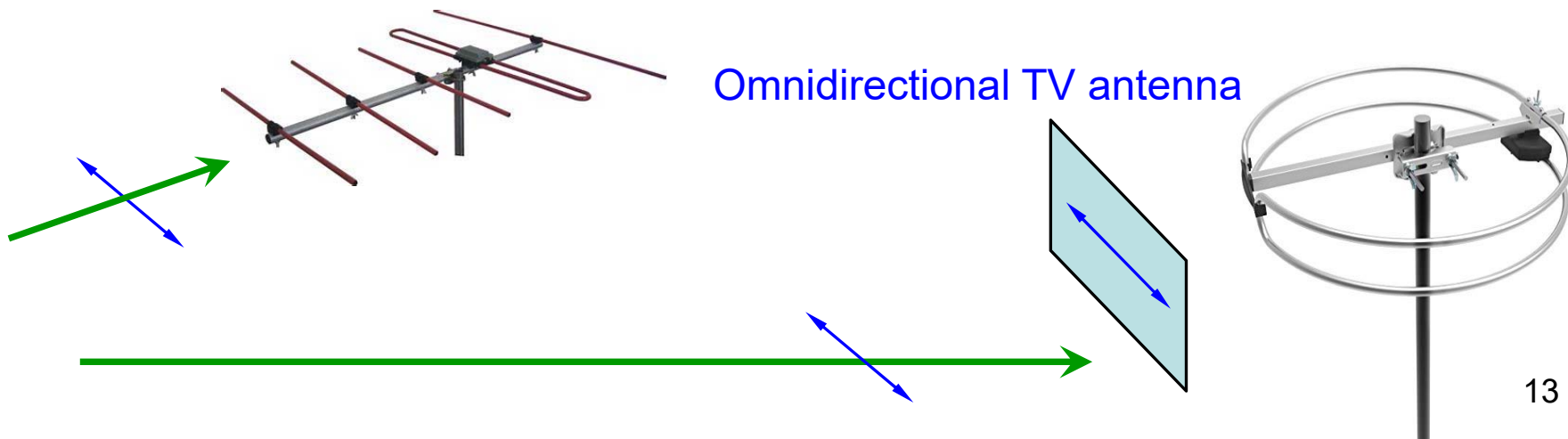
Inertial Confinement Fusion is a promising option to provide massive, clean, and affordable energy for mankind in the future. The present status of research and development is hindered by hydrodynamical instabilities occurring at the intense compression of the target fuel by energetic laser beams. A recent patent combines advances in two fields: Detonations in relativistic fluid dynamics (RFD) and radiative energy deposition by plasmonic nano-shells. The initial compression of the target pellet can be decreased, not to reach the Rayleigh–Taylor or other instabilities, and rapid volume ignition can be achieved by a final and more energetic laser pulse, which can be as short as the penetration time of the light across the pellet. The reflectivity of the target can be made negligible as in the present direct drive and indirect drive experiments, and the absorptivity can be increased by one or two orders of magnitude by plasmonic nano-shells embedded in the target fuel. Thus, higher ignition temperature and radiation dominated dynamics can be achieved with the limited initial compression. Here, we propose that a short final light pulse can heat the target so that most of the interior will reach the ignition temperature simultaneously based on the results of RFD. This makes the development of any kind of instability impossible, which would prevent complete ignition of the target.

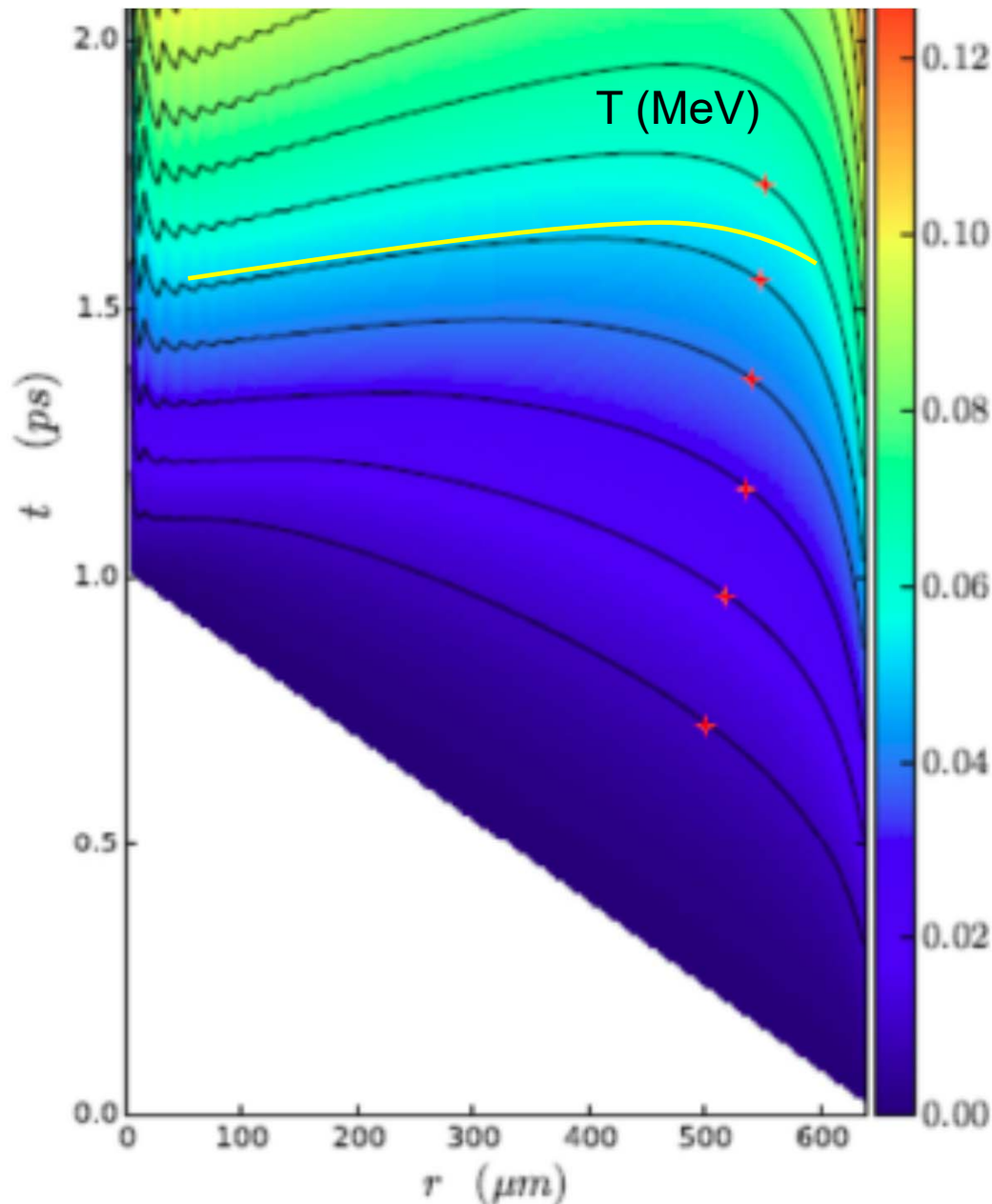
Metal nanoparticles (MNP) and their optical properties



[Martin Greve, IFT Seminar, Fall (2017) for PV Solar panels]

Localized Surface Plasmon Resonance (LSPR)!





The absorption coefficient is **linearly** changing with the radius: In the center, $r = 0$, $\alpha_K = 30 \text{ cm}^{-1}$ while at the outside edge $\alpha_K = 8 \text{ cm}^{-1}$.

The temperature is measured in units of $T_1 = 272 \text{ keV}$, and $T_n = n T_1$.

Simultaneous, volume ignition is up to 0.9 R, so 73% of the fuel target!

Problem:

Not easy to realize, 4π irradiation geometry → **80-192 laser beams**

In the earlier estimates we did establish that:

- The principle of simultaneous volume ignition is theoretically possible with nano-spheres or nano-rods (i.e. nano-antennas).
 - We **did not** estimate the necessary laser energy need, and assuming extremely large laser input energy we neglected the losses.
-
- However the basic principle, the **simultaneous all volume ignition** can be achieved in 3D, 2D and 1D geometry also !!!
(Einstein's synchronization of watches.)
 - We did calculate the realistic energy balance
 - → So, we repeated realistic estimates in 1D with a flat (coin or rod like) target.

Thick Coin like target - New Developments

Thickness of the target is: h

h depends on pulse energy, ignition energy, target mass, ...

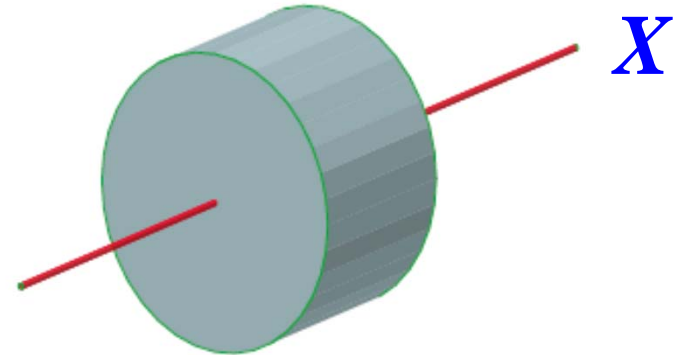


Figure 1: (color online) The target still should be compact to minimize the surface effects. The irradiation is performed along the x -axis from both sides towards the target. The laser beam should be uniform hitting the whole face of the coin shaped target.

Radiation dominated implosion with flat target

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¹

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² Dept. of Optics and Quantum Electronics, Univ. of Szeged, Hungary

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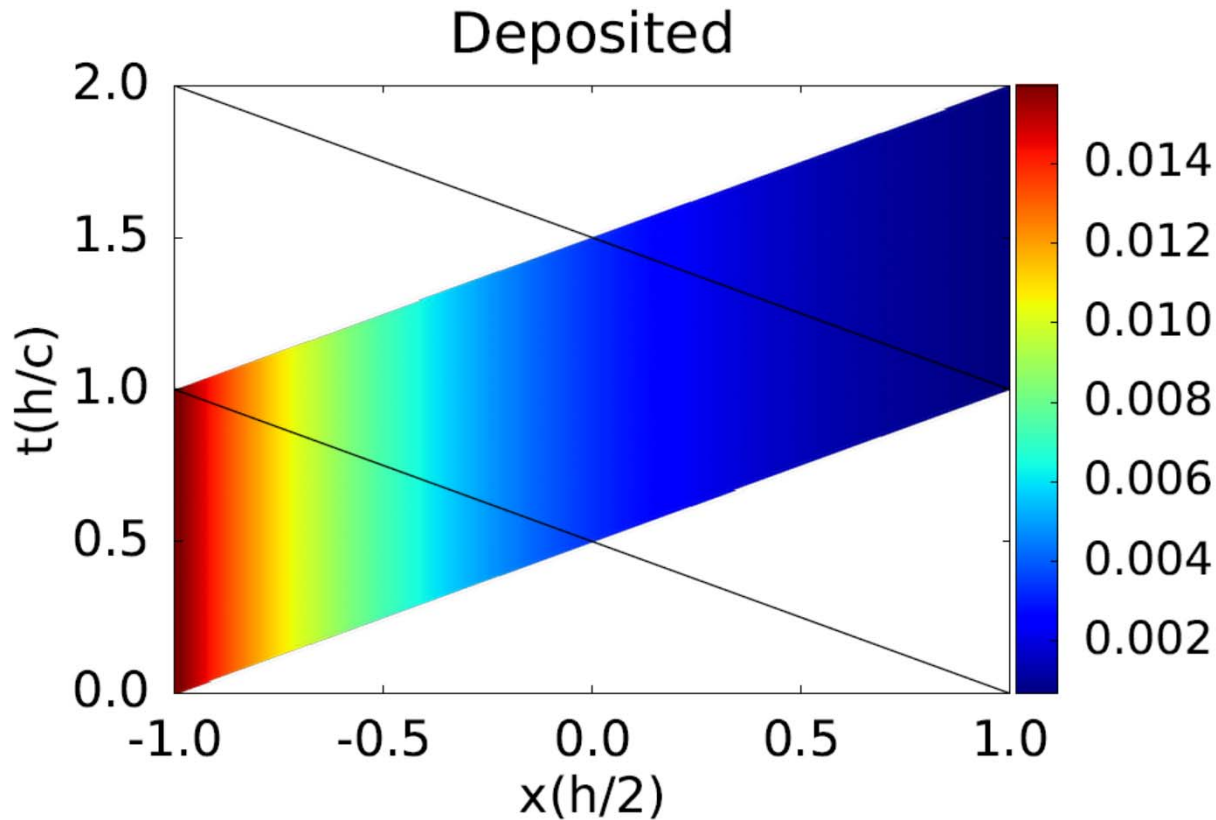
⁵ Hungarian Academy of Sciences, Budapest, Hungary

⁶ Wigner Research Centre for Physics, Budapest, Hungary

⁷ National Research Center "Kurchatov Institute" Moscow, Russia

[arXiv:1903.10896v3

- in press, PWP]



**Without nano
antennas**

The deposited energy from laser irradiation from one side only. The absorption is constant, this leads to an exponentially decreasing energy deposition, and only a negligibly small energy reaches the opposite end of the target.

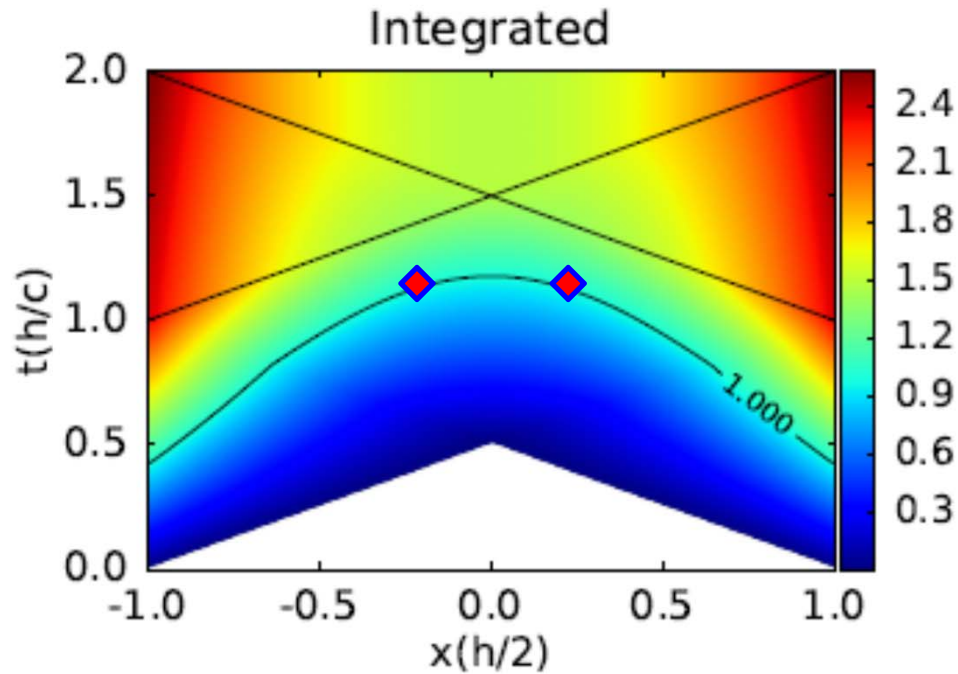


Figure 4: (color online) Integrated energy up to a given time in the space-time across the depth, h , of the flat target. The color code indicates the temperature, T , reached in a given space-time point, in units of the critical temperature, (T_c). The contour line $T = 1$, indicates the critical temperature, T_c , where the phase transition or the ignition in the target is reached. This contour line, compared to the one in Fig. 3, is never constant in time, indicating no simultaneous whole volume transition or ignition. The time-like (causally unconnected) part of the transition takes place only in the central $\sim 15\%$ of the target volume. The two straight lines indicate the light-cones originating at the outside edges of the target at the ending of the irradiation pulse.

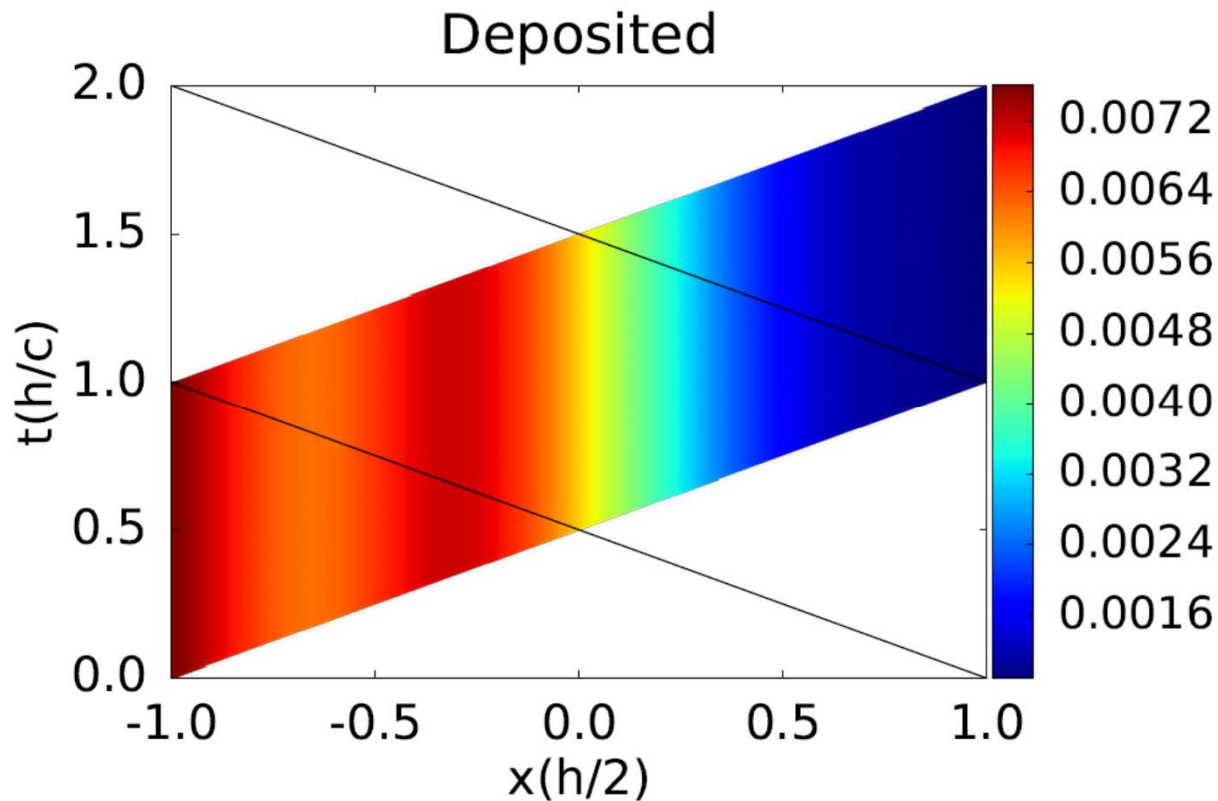
Without nano antennas

Irradiation from both sides.

Exponential decrease of deposited energy. Due to the already deposited energy, less energy reaches the middle →

The front and back surface is heated up but the middle is not!

Pulse length is: $t_p = h/c$



With nano antennas

The absorptivity is increased towards the center, due to the implanted nano antennas.

The deposited energy from laser irradiation from one side only. The absorption is modified by nano antennas so that the absorptivity is increasing towards the middle, so that the deposited energy is constant up to the middle. Then the absorptivity is decreasing, but hardly any energy is left in the irradiation front. Thus again only a negligibly small energy reaches the opposite end of the target.

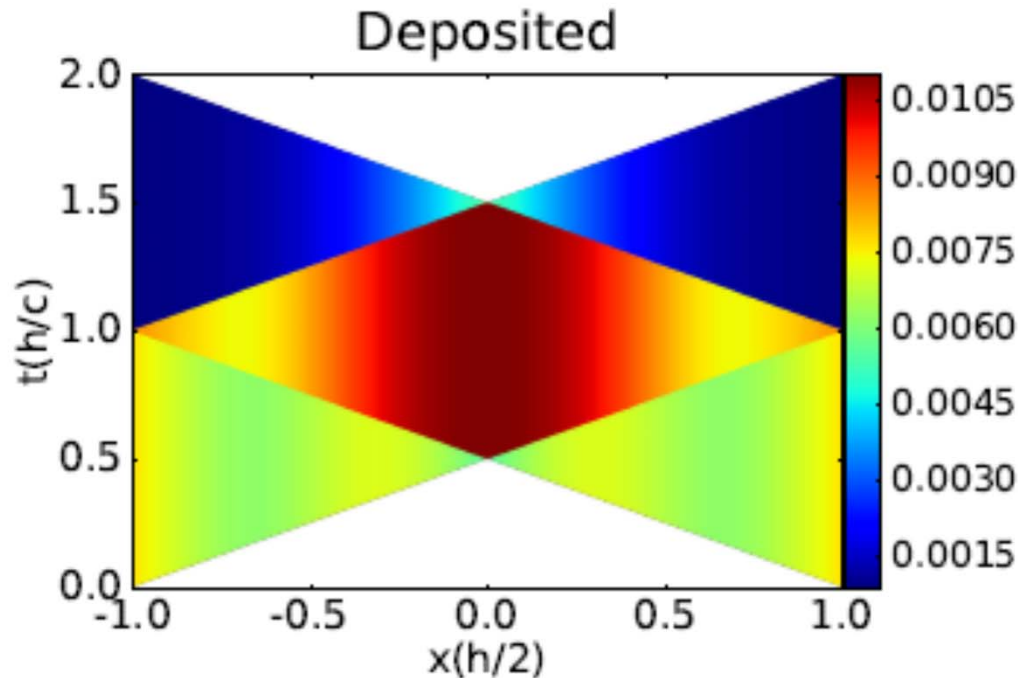


Figure 2: (color online) Deposited energy per unit time in the space-time across the depth, h , of the flat target. The time is measured in units of (h/c) , where c is the speed of light in the material of the target. The irradiation lasts for a period of $\Delta t = h/c$ the time needed to cross the target. The irradiated energy during this time period is Q from one side, so it is $2Q$ from both sides together.

The color code indicates the deposited energy per unit time and unit cross section (a.u.). The deposited length is $\Delta x = c\Delta t$. Note! The absorptivity in this case $\alpha_K \neq \text{const}$. For more details please see Appendix B.

With nano antennas

Irradiation from both sides.

Ignition energy is: Q_i/m
 e.g. for DT target: $Q_i/m = 27 \text{ kJ/g}$

→ if we have $Q = 100 \text{ J}$, then we can have a target mass:

$$m_{DT} = Q / Q_i \text{ g} = 3.703 \text{ mg.}$$

Then with m_{DT} and ρ_{DT} given we get the DT-target's volume, V_{DT} and $h_{DT} = 2.67 \text{ mm}$.

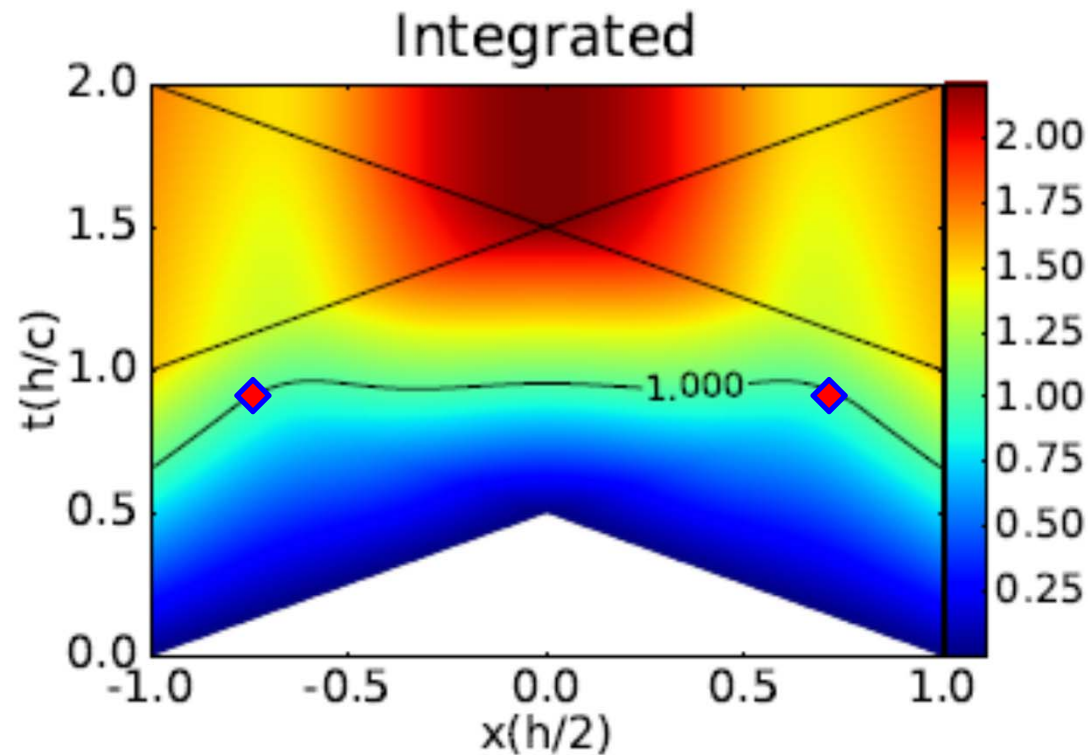


Figure 3: (color online) Integrated energy up to a given time in the space-time across the depth, h , of the flat target. The color code indicates the temperature, T , reached in a given space-time point, in units of the critical temperature, (T_c) . The contour line $T = 1$, indicates the critical temperature, T_c where the phase transition or the ignition in the target is reached. This contour line is almost at a constant time, indicating simultaneous whole volume transition or ignition. The irradiated energy, Q is chosen so that, $1Q$ irradiation will achieve the critical temperature.

With nano antennas

Ignition is reached at contour line $Q = 1$.

[L. P. Csernai, M. Csete, I. N. Mishustin, A. Motornenko, I. Papp, L. M. Satarov, H. Stöcker, N. Kroo, *arXiv:1903.10896*, *Submitted to MRE*]

Test of principles with smaller pulse energy

- 1 Relativistic time-like (simultaneous) ignition
- 2 Using nano antennas to reach whole volume uniform ignition or transition
- 3 Using 1D geometry, with two beams from opposite direction for realizability

Let us take

a $P = 30 \text{ mJ}$, 1 ps laser,

a polylactic acid (PLA) target with $T = 150 \text{ C}$ melting temperature, $Q_i/m = 28 \text{ J/g}$

this leads to $m_i = 0.32 \text{ mg}$ target mass, and $h_i = 683 \text{ }\mu\text{m}$ target thickness.

The melting transition profile can be checked by simple and affordable analysis by microscope.

The distribution, and absorption properties of implanted nano antennas can be well optimized.

European Laser Infrastructure – Szeged, HU



ELI-ALPS Szeged:
EU Extr. Light Infrastructure
Attosec. Light Pulse Source

2PW High Field laser
10 Hz, $<10\text{fs}$, **20 J**



HAS Wigner RCP, Budapest

Gagik P. Dzsotjan, József Bakos, Gábor Demeter,
Dávid Dzsotjan, Miklós Kedves, Béla Ráczkevi,
Zsuzsanna Sörlei, Péter Lévai



Ti-Si Hydra L. 30mJ 10Hz 40fs

Laser wake acceleration of protons for radiation therapy

- proton beam energy is deposited at a location of a certain depth [Bragg peak]
- tumor treatment with minimal side damage (compared to other radiation therapies)
- target is low density (~ like water or more)
- Collaboration with Peking University, China

These features are similar to the needs of laser induced ICF with nano-plasmonics!
Deposition at a depth via the Bragg peak is an alternative way to get volume ignition

Gábor Veres, István B. Földes, Márk Aladi, Imre Ferenc Barna, Róbert Bolla,
Zsolt Kovács, Mihály Pocsai, Dániel Dunai, Gábor Anda et al.

Fusion plasma diagnostics, ITER, JET etc.

Péter Dombi, Péter Rácz, Norbert Kroo et al.

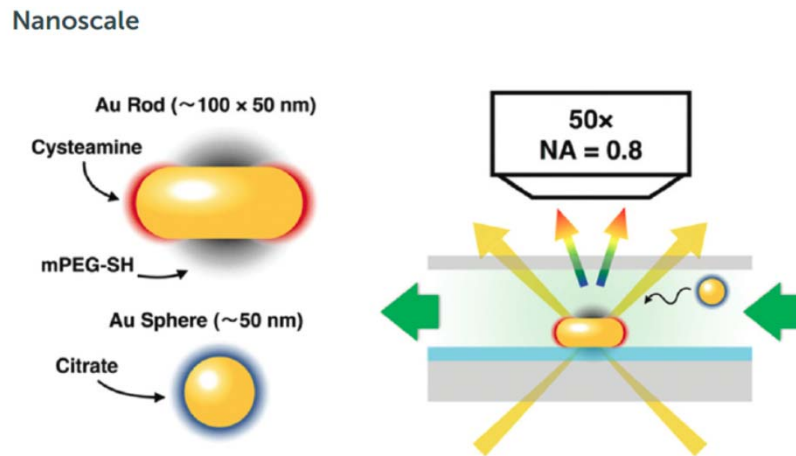
Laser induced nano-plasmonics

HAS Centre for Energy Research, Budapest Inst. for Technical Physics and Materials Science

András Deák, S. Pothorsky, D. Zámbo, D. Szekrényes, Z. Hajnal, Béla Pácz et al.

Nano-particle assembly at the single particle level

- manufacturing of Au nano-shells and nano-rods
- imbedded in different concentrations in carriers
- polarized target constructions with nano-rods (for polarized laser irradiation)
- testing resonant light absorption



Scheme 1 Schematics of the prepared nanoparticles (left) and the measurement arrangement (right). The patchy nanorods are first immobilized on ITO covered substrates, then the aqueous nanosphere solution is introduced and changes in the scattered spectrum upon binding detected.

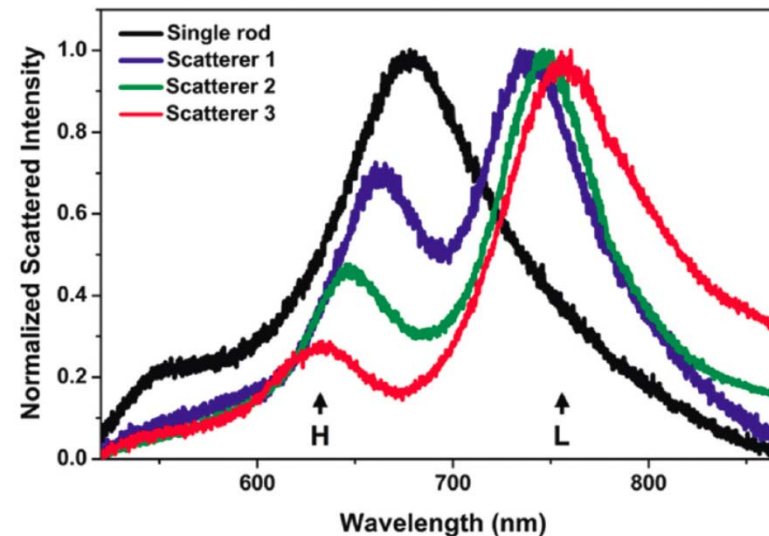
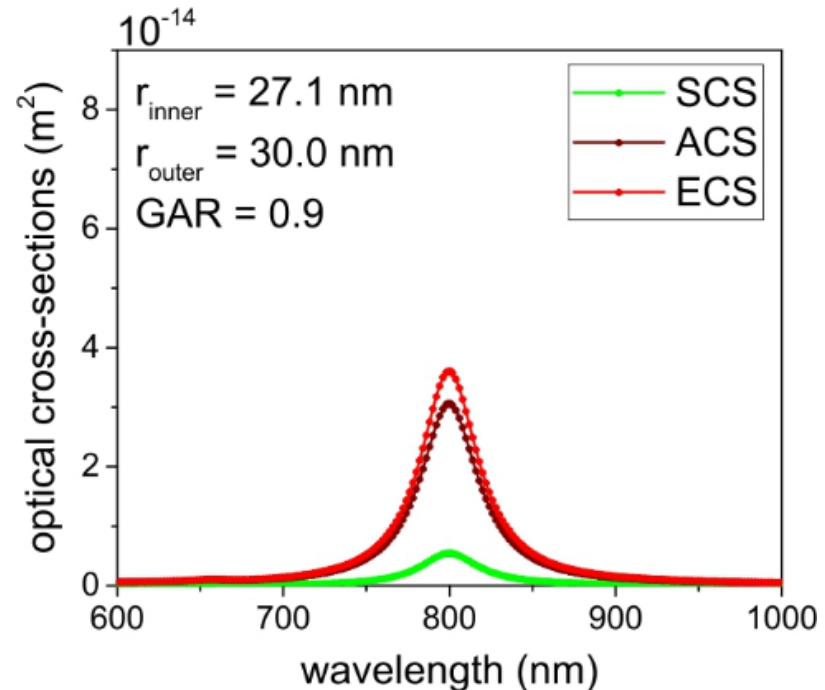
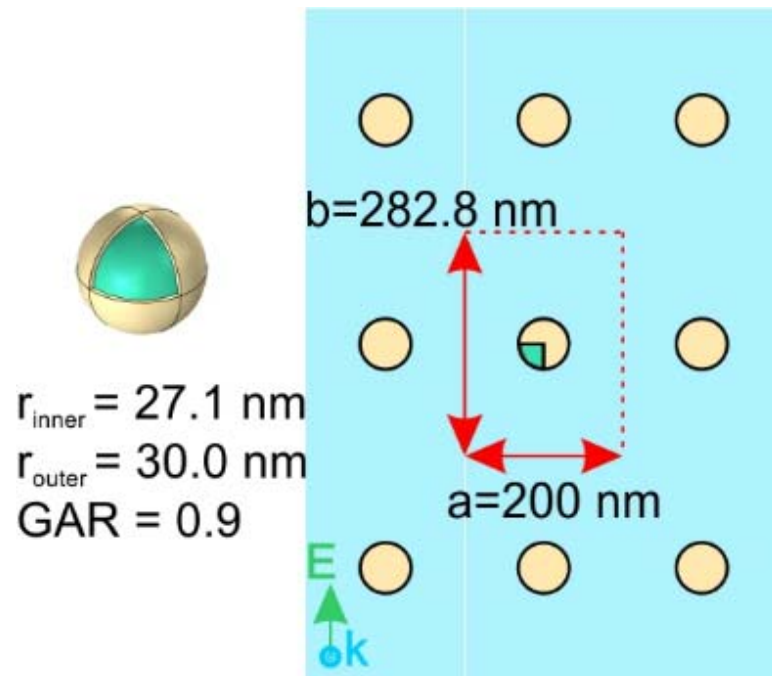


Fig. 1 Typical scattering spectra measured *in situ* in the liquid cell. The appearance of the high (*H*) and low (*L*) energy peaks surrounding the dip indicate the formation of a heterodimer. The spectrum of a single nanorod before the assembly is shown for reference.

Nano-particle absorption

The target absorptivity is increased via core-shell type plasmonic nano-shells. Calculations via solving the Maxwell equations, and evaluating the ohmic heating were performed.



1 ps laser pulse length, $\lambda = 800 \text{ nm}$, One-sided & two-sided irradiation tested, 85-100 % absorption in the target length h . Nano-antenna shapes, layer configurations, layer distribution varied & analyzed.

[M. Csete, et al., University of Szeged, HU]

50 Jahre GSI



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GSI ist Mitglied bei

■ Laser Parameter: PHELIX

	Langer Puls	Kurzer Puls
Pulslänge:	1 - 10 ns	0.5 - 20 ps
Energie:	0.3 - 1 kJ	bis 200 J
Max. Intensität:	10^{16} W/cm ²	$2 \cdot 10^{21}$ W/cm ²
Kontrast:	50 dB	bis 120 dB

■ Wissenschaftliche Veröffentlichungen

PHELIX wird genutzt, um verschiedene wissenschaftliche Gebiete, meist in der Plasma- und Atomphysik, zu erforschen. Auf der Seite mit den [🔗 PHELIX Publikationen](#), erhalten Sie zusätzliche Informationen, sowie die aktuellsten wissenschaftlichen Veröffentlichungen, die mit den Daten aus Experimenten mit dem PHELIX-Laser erarbeitet wurden.

<https://www.gsi.de/work/forschung/appamml/plasmaphysikphelix/phelix.htm>

Available resources:

	Pulse energy	Pulse frequency	Pulse length
LLNL NIF 192 laser 3D (1 laser:	11.2 kJ	1/day	~ 10-30 ns
ELI-APLS 2PW High field laser	20 J (34 J)	10 Hz	< 10 fs
BELLA, L. Berkeley Nat. Lab.	34 J	1 Hz	< 30 fs
GSI/FAIR PHELIX laser	200 J	1/90min	> 0.5 ps
Wigner - Coherent Ti-Si Hidra L. (upto 100 mJ)	30 mJ	10 Hz	40 fs
Optimal for laser induced ICF with nano-plasmonics tests	~ 100 J	1 Hz	1-10 ps !

Experimental test of similar configuration @ ShenGuang-II Up, Shanghai :

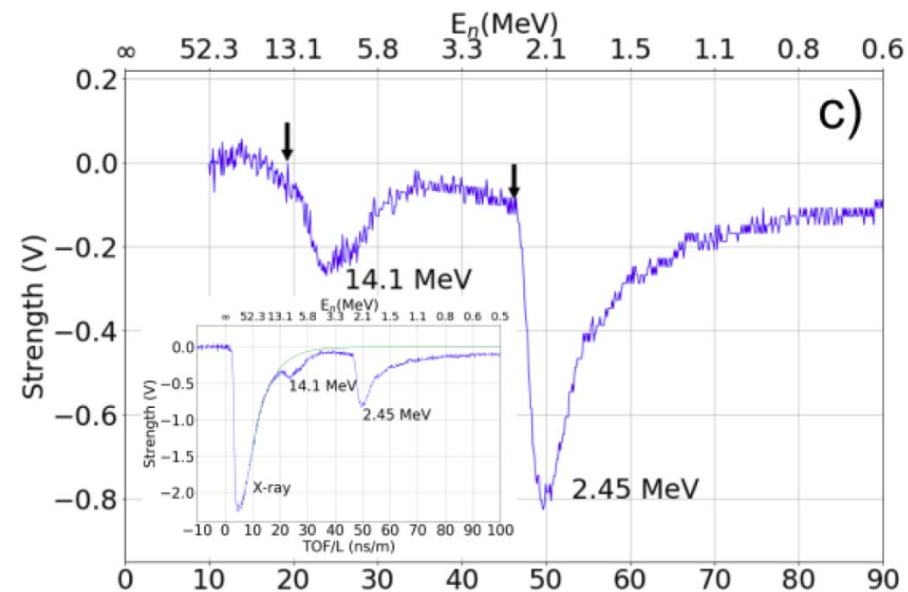
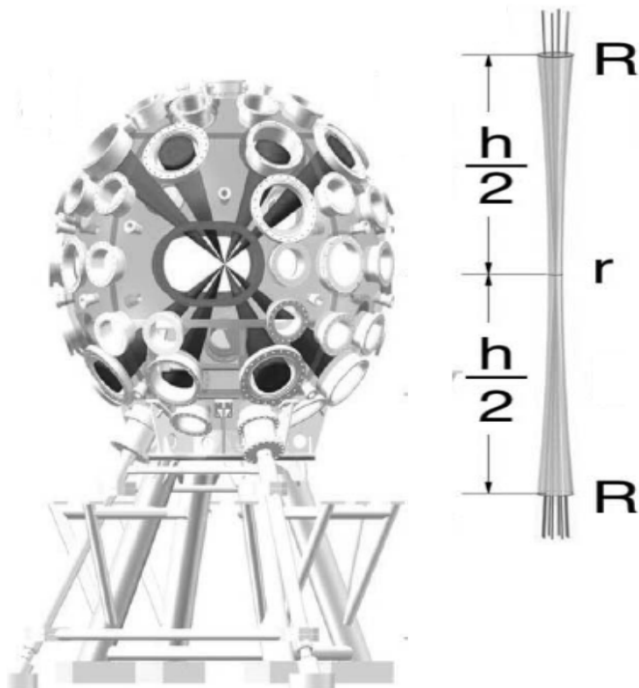
Nuclear probes of an out-of-equilibrium plasma at the highest compression
Phys. Lett. A 383 (2019) 2285-2289.

G. Zhang^{a,b,*}, M. Huang^c, **A. Bonasera^{d,e,*}**, Y.G. Ma^{f,b,i,*}, B.F. Shen^{g,h,*}, H.W. Wang^{a,b},
 W.P. Wang^g, J.C. Xu^g, G.T. Fan^{a,b}, H.J. Fu^b, H. Xue^b, H. Zheng^j, L.X. Liu^{a,b}, S. Zhang^c,
 W.J. Li^b, X.G. Cao^{a,b}, X.G. Deng^b, X.Y. Li^b, Y.C. Liu^b, Y. Yu^g, Y. Zhang^b, C.B. Fu^k,
 X.P. Zhang^k

4 (up) + 4(down) lasers
 Target thickness, h ($3.6\mu\text{m}$ - 1mm)
 & radius, R , (150 - $400\mu\text{m}$) were
 varied.

Total pulse energy 1.2kJ (2ns) for 8 beams.
 Shortest (250ps) pulses -> 100s MeV ions >
 non-thermal distr. = directed ion acceleration

Typical fusion neutron energies were measured
 & used to extract the target density.



Experimental test of similar configuration @ ShenGuang-II Up, Shanghai :

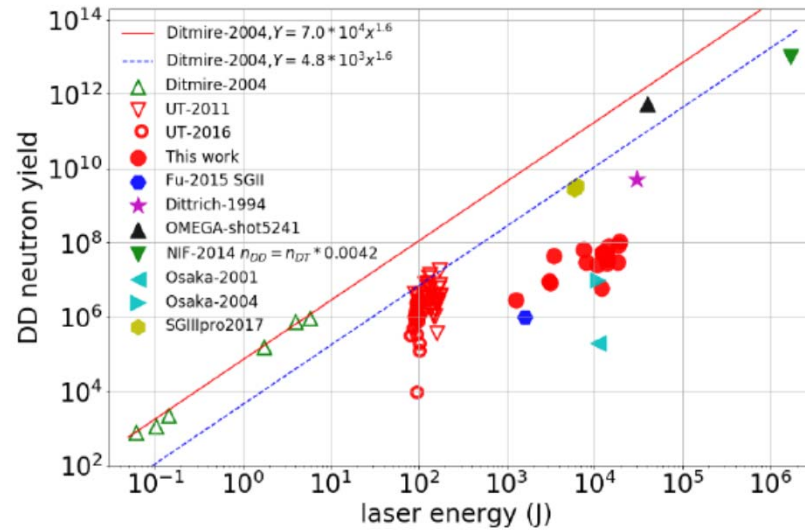


Figure 3: (color online) Fusion yield as function of laser energy. Different experimental results Ditmire-2004[40], UT-2011[20], UT-2016[19], Fu-2015 SGII[45], Dittrich-1994[49], NIF-2014[48], Osaka-2001[46], Osaka-2004[47], OMEGA-shot5241[41] and SGIIpro2017[42] are indicated in the inset.

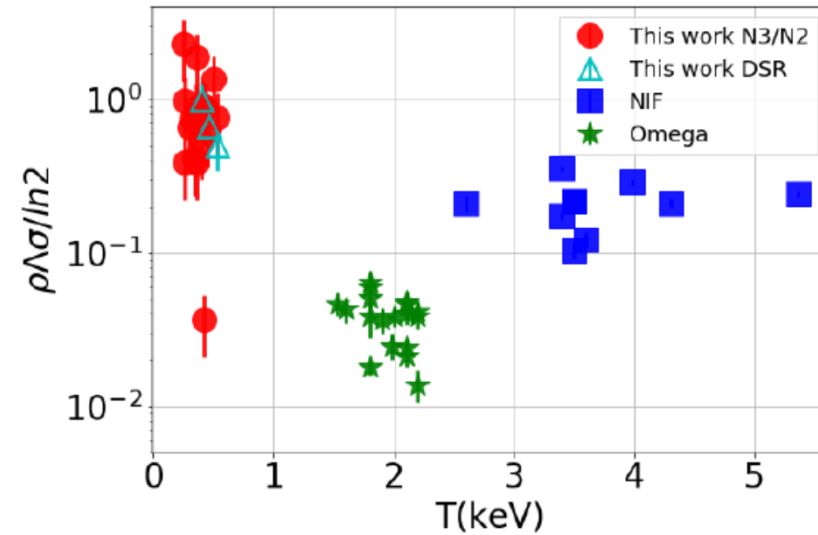
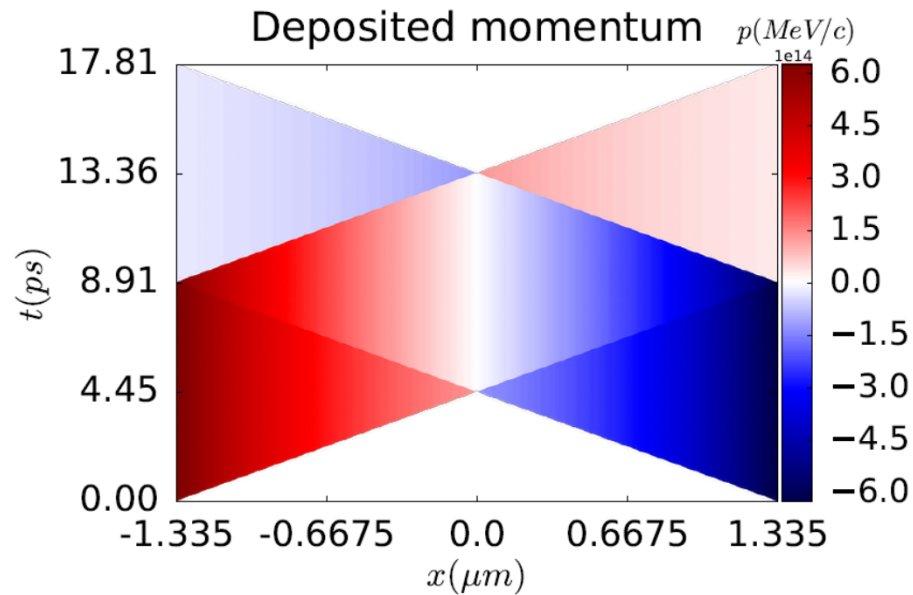


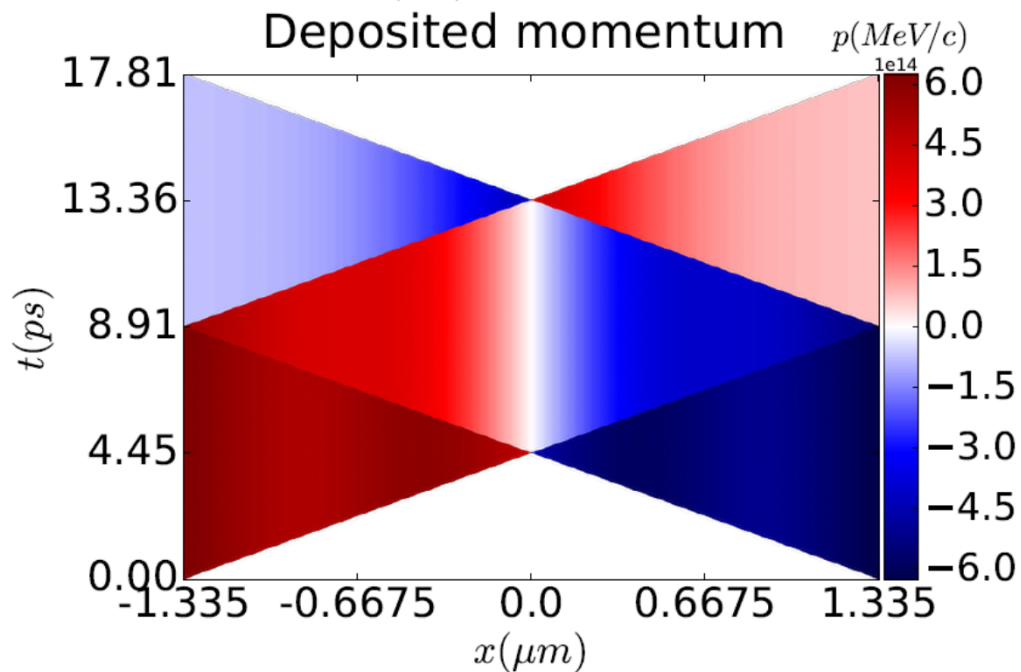
Figure 4: (color online) $\Lambda \rho \sigma / \ln 2$ obtained from eq.(4) vs T from eq.(1). Omega and NIF data are derived from the experiments[25], using the Down Scatter Ratio[23, 21]. Our results using the DSR method (N_4/N_3) are given by the open triangle symbols in good agreement with the N_3/N_2 ratios.

Stimulated by these considerations we decided not to fight non-equilibrium effects but rather enhance them, i.e. study plasmas highly compressed and completely out of equilibrium.

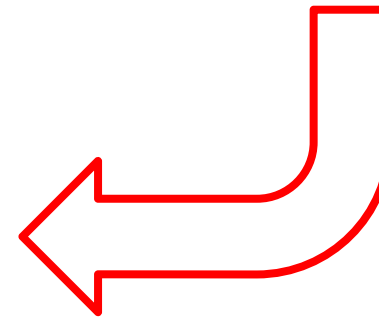
Deposited momentum



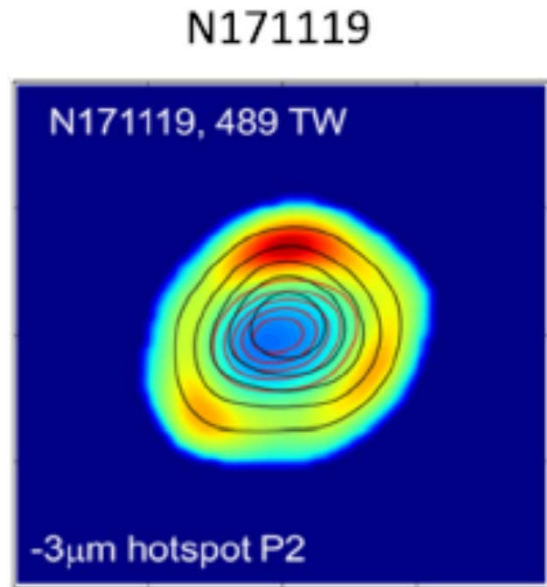
Leads to compression w/o external ablator surface
[Gyulassy, Csernai, NPA 1986]
→ leads to QGP formation w/o RT instabilities



Implanted nano-antennas increase further the compression !!

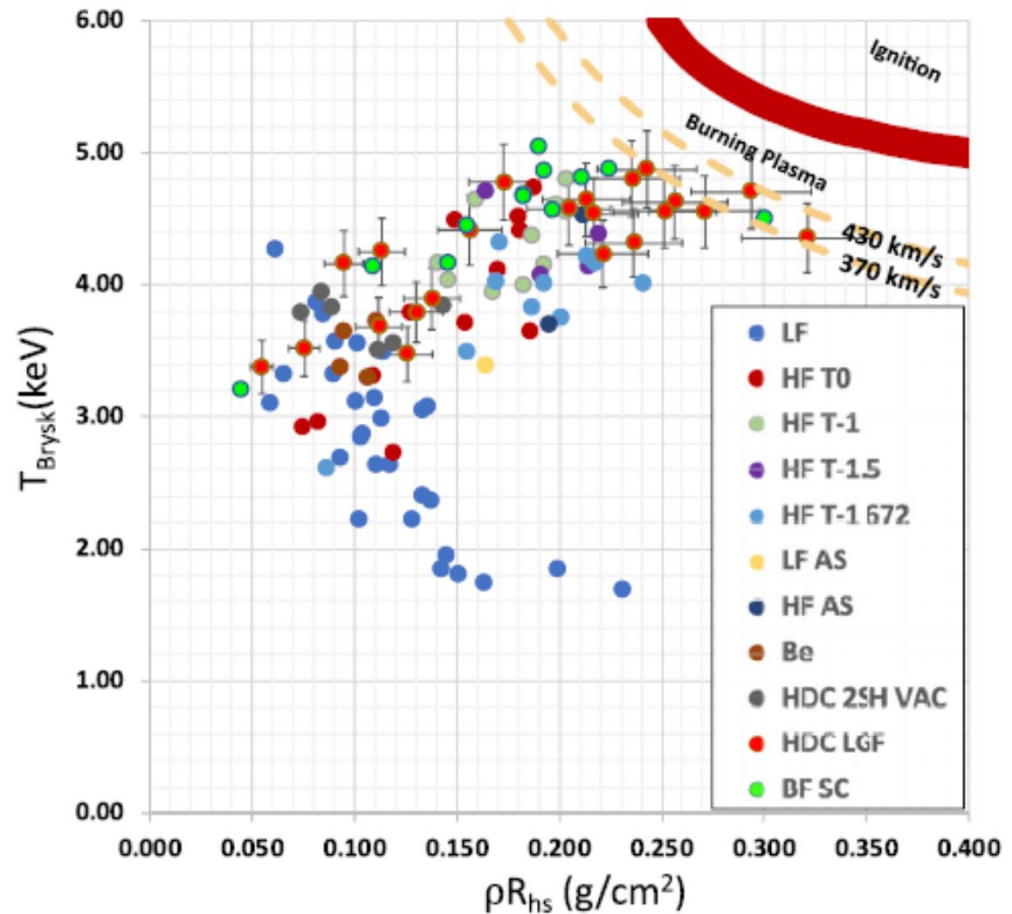


LLNL – NIF April 2019



Hot-spot, 3 μ m, burning spreads by α -heating: boundary is approached.

[O.A. Hurricane et al., Phys. Plasmas 26, 052704 (2019) April.]



Thus, ultra-relativistic heavy ion physics lead to discovery **Quark Gluon Plasma (QGP)**, but also to advances in

(i) relativistic fluid dynamics (RFD).

With **(ii) nano technology** this may revolutionize in a simple, and **(iii) affordable 1D geometry** the technological development of

(→ iv) Inertial Confinement Fusion.

This methodology may be used to study other dynamical, radiation dominated processes.



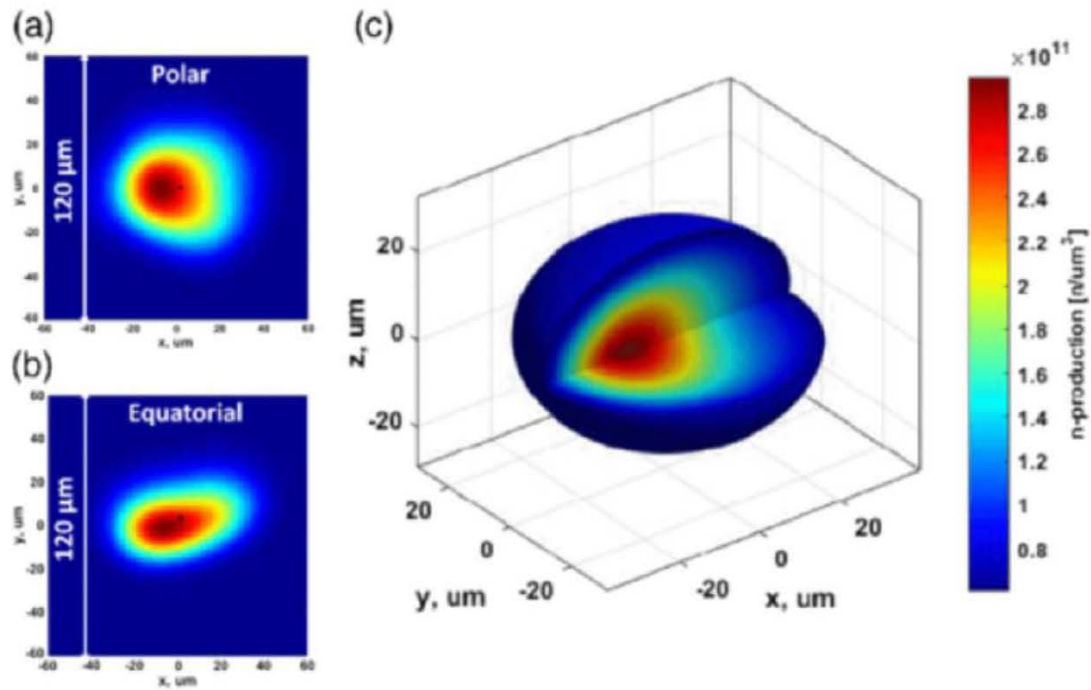


Figure 3

Shot N170601. (a) Polar neutron image. (b) Equatorial neutron image. (c) Three-dimensional reconstructed neutron volume of the hot spot.

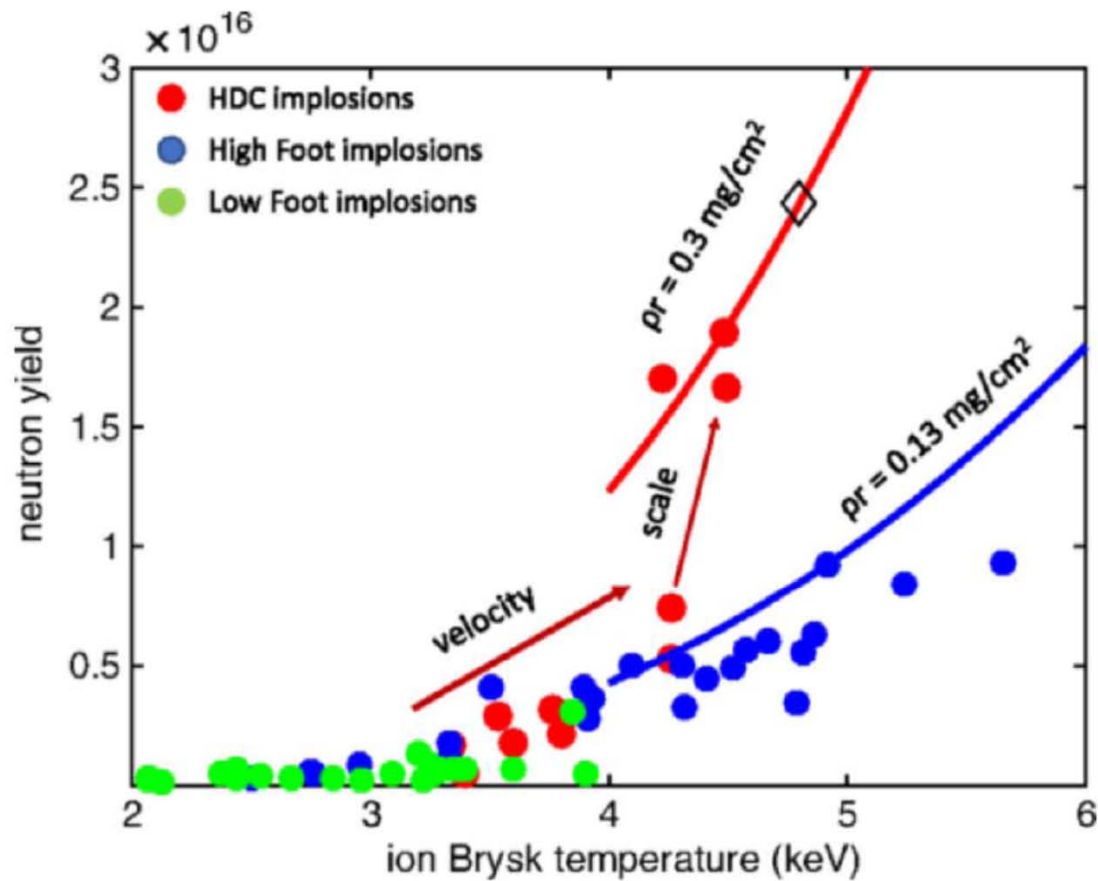


Figure 4

Total DT neutron yield as a function of ion temperature, red dots are doped HDC implosions, blue dots are high foot implosions, green dots are low foot implosions. The neutron yield is plotted against the lowest burn averaged DT ion temperature measured by NTOF detectors (Brysk temperature). For high foot implosions, the Brysk temperature is estimated to be up to a keV higher due to flows in the hot spot. Black diamond is the point where α deposited energy equals bremsstrahlung and conduction losses. Solid curves are a yield extrapolation with temperature using a constant ρr and adiabat.

